Spontaneous Symmetry Breaking of a Hinged Flapping Filament Generates Lift

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Advantage/disadvantage of *passive flexible* parts/appendages in animal propulsion?















Re >> 1, birds/fish use *reciprocal flapping* for lift and/or propulsion







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Heaving, rigid, symmetric foil (thrust)



(inverted von Karman vortex street)





Re >> 1, birds/fish use *reciprocal flapping* for lift and/or propulsion



Heaving, rigid, cambered foil (plus lift)





Re >> 1, birds/fish use *reciprocal flapping* for lift and/or propulsion



Heaving, *flexible*, cambered foil (plus ...)

J. Guerrero, PhD Thesis, 2009, UNIGE



Some animals live between these worlds: $Re \approx 10 - 100$ (and present flexible appendages)



Clione Antarctica ("sea slug"), a marine mollusc which switches mobility strategy with adulthood:

rowing cilia \rightarrow flapping wings

Childress & Dudley, JFM 2004



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Dauptain, Favier & Bottaro, JFS 2008



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fly



butterfly



mosquito



"pop-up" feathers

In Nature "rough" is rule, not exception ...



Is there some dramatic change in the interaction between a "free" body and a fluid as Re /

What is the role of apparently *inert* filaments/cilia/tentacles/scales/feathers?





Symmetry breaking!

Filament flap asymmetrically:

- induced force and torque
- reduced drag





$$\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} = -\nabla p + \frac{1}{Re} \nabla^2 \mathbf{u}$$
$$\nabla \cdot \mathbf{u} = 0$$

No-slip

 $\mathbf{u}=0\qquad\text{on}\quad\Gamma$





$$\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} = -\nabla p + \frac{1}{Re} \nabla^2 \mathbf{u} + \int_{\Gamma} \mathbf{f}(\zeta) \delta(\mathbf{x} - \zeta) \, \mathrm{d}\zeta$$
$$\nabla \cdot \mathbf{u} = 0$$

No-slip

 $\mathbf{u} = 0$ on Γ

Flow field: Eulerian grid Boundary: Lagrangian points Boundary force to enforce no-slip condition Projection method (Taira & Colonius, JCP 2005)





$$\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} = -\nabla p + \frac{1}{Re} \nabla^2 \mathbf{u} + \int_{\Gamma} \mathbf{f}(\zeta) \delta(\mathbf{x} - \zeta) \, d\zeta$$

$$\nabla \cdot \mathbf{u} = 0$$

No-slip

$$\mathbf{u}(\Gamma) = \dot{\zeta} \qquad \zeta = \zeta(s,t) \quad s: \text{ arclength}$$

Filament dynamics : Euler-Bernoulli



Peskin, *Acta Numerica* 2002 Kim & Peskin, *PoF* 2007





The cylinder alone: symmetric wake

– Reynolds number
$$Re=rac{{f U}D
ho_f}{\mu}$$

– Vortex shedding for $Re > Re_c$ with frequency f_c



(classical von Karman vortex street)





The filament alone: symmetric flapping



The filament alone (soap film experiments)



Zhang, Childress, Libchaber & Shelley, *Nature* 2000 Shelley & Zhang, *ARFM* 2011



The filament alone (theory)



FIG. 2 (color online). Snapshots of the flag for fixed mass $(R_1 = 0.3)$ and decreasing rigidities R_2 . (a) The observed flapping mode at $R_2 = 0.01445$, about a factor of 2 below the critical $R_2 = 0.0262$; (b) a higher energy flapping mode for $R_2 = 0.0138$; (c) a chaotic flapping mode at $R_2 = 0.0025$. (d) Comparison of experimental flag snapshots in [18], Fig. 12b, with model shapes. In both cases $R_1 = 0.37$ (see text for R_2).

Inviscid flow theory: vortex sheet (filament) + shed free vortex sheet \rightarrow Biot-Savart integral (for velocity \perp to the filament) + Euler equation

Alben & Shelley, PRL 2008



The filament alone (numerical simulations)

• Flapping when $Re > 10^3$ $R_1 > 0$ $R_2 < R_{2,c}$





Cylinder + filament: something happens ...

$$L = \frac{L_s}{D}, \quad Re = \frac{UD}{\nu} \quad R_1 = \frac{\rho_s}{\rho_f D} \quad R_2 = \frac{B}{\rho_f U^2 D^3},$$





$Re = 100, R_1 = 0.1, R_2 = 0.005, L = 3$ and L = 1.5



time = 315.05 L = 1.50





Bagheri, Mazzino & Bottaro, PRL 2012





L	R_2	C_d	C_l	C_q	f_c
0.0	-	1.36 ± 0.01	0.00 ± 0.34	0.00	0.164
3.0	0.005	1.28 ± 0.06	0.00 ± 0.23	0.00	0.157
1.5	0.005	1.32 ± 0.08	0.18 ± 0.28	0.01	0.159
1.5	0.100	1.23 ± 0.05	0.21 ± 0.24	0.02	0.145

increasing $R_2 \rightarrow$ increased rigidity of the structure





FIG. 4 (color online). Left, center, and right columns correspond, respectively, to long/flexible (case ①), short/flexible (case ②) and short/stiff filaments (case ③). Top row shows instantaneous snapshots of the maximum value of the FTLE and filament position (solid line). Bottom row shows mean filament position (solid line) and mean pressure distortion. The latter is defined as the difference between the time-averaged pressure field with and without the filament. Positive (negative) distortions are plotted with solid black (gray) contour lines; positive pressure distortion in the leeward side of the cylinder signals reduced pressure drag, while asymmetric distortion with respect to the x axis indicates that a net y force is generated.



Choice of observable







A symmetry-breaking bifurcation occurs when vortices and structures resonate ...



Beam equation

Equation governing unforced beam

$$R_1 Y_{tt} + R_2 Y_{ssss} = 0$$

• Eigenfrequency

$$f_s = \sqrt{\frac{R_2}{R_1 L^4}}$$



Resonance condition

- Free vibrations of filament f_s
- Vortex shedding frequency f_c
- If $f_s \ll f_c$ filament with slow reaction time
- If $f_s \gg f_c$ filament reacts instantaneously
- Thus $f_s \sim f_c$ separates two different regimes

$$L_r = \left(\frac{R_2}{R_1 f_c^2}\right)^{1/4}$$



- Energy
$$E = \frac{1}{2} \int_0^L R_1 |\dot{\zeta}|^2 + R_2 |\zeta_{ss}|^2 ds$$

- Rescaled with filament density and length

$$(\rho_f, D) \to (\rho_s, L_s)$$

Flapping synchronized with vortex shedding, time scale

 \rightarrow rescaled non-dimensional filament energy

$$\tilde{E} = \frac{R_1}{L^3} E$$



Resonance



• Resonance: L = 1.25 (flexible) L = 2.25 (rigid)



Resonance

	Resonance (theoretical)	Resonance (computed)	Bifurcation (computed)
Flexible	1.25	1.25	1.6
Rigid	2.6	2.25	2.25



Can filaments increase drift?

- Efficient wind-borne seed dispersal
 - Side force due to symmetry breaking may increase drift







(Burrows, New Phytol. 1975)

Can filaments reduce drag "optimally" (because of compliance) as opposed, i.e., to the (sub-optimal) pressure drag reduction of golf balls?



How to model a passive, compliant hairy/feathery coating?



sea otter

egret



GOAL: instead of a single flexible flap, let's model a continuous *hairy/feathery* coating to affect lift and drag



Numerical challenges





- Model mechanical properties of biological surfaces
- Structures with large displacements and large rotations
- Interaction between multiple structures

Coupling between a layer of oscillating densely packed structures and a unsteady separated boundary layer



The initial configuration



Circular cylinder, Re=200

Model of the layer?

Porous, anisotropic and compliant







Aerodynamic performances

	<c<sub>d></c<sub>	C _d '	C _l '	St
Case 1	1.3689 (1.39;1.356)	0.0274	0.4381	0.199 (0.199;0.198)
Case 2	3.1464	0.1943	1.1376	0.1946
Case 3	1.3035	0.0207	0.3839	0.1916
Case 4	1.2109	0.012	0.3008	0.1661

(Bergmann et al., PoF 2005; He et al., JFM 2000)



Aerodynamic performances

		<c<sub>d></c<sub>	C _d '	C _l '	St
Case 1	\bigcirc	ref	ref	ref	ref
Case 2	Y X	+130%	+608%	+160%	-2.21%
Case 3		-4.78%	-24.54%	-12.37%	-3.71%
Case 4	<u>}</u>	-11.54%	-56.09%	-31.34%	-16.53%



Physical mechanism



inerence of time-averaged pressure ne

<P with hair>-<P ref>

Physical mechanism



Contours of vertical velocity

Movements of reference cilia



Contours of vertical velocity

Force field

The hairy layer counteracts flow separation





Optimal self-adaptive hairy layer

15% drag reduction

40% reduction in lift fluctuations



Favier, Dauptain, Basso & Bottaro, JFM 2009





Consider a airfoil: the control elements (the "feathers") must be placed in the position of largest *sensitivity* to achieve an effect



Hairfoil



Consider a hairfoil: the control elements (the "feathers") must be placed in the position of largest sensitivity to achieve an effect



Hairfoil (simulations)

. . .



Potential applications:

MAV/UAV Wind turbines Hydraulic machines (cavitation?) Sound mitigation

Venkataraman & Bottaro, PoF 2012



Hairfoil (experiments)





Kunze & Brücker, CRAS 2012

... but this is another story!







