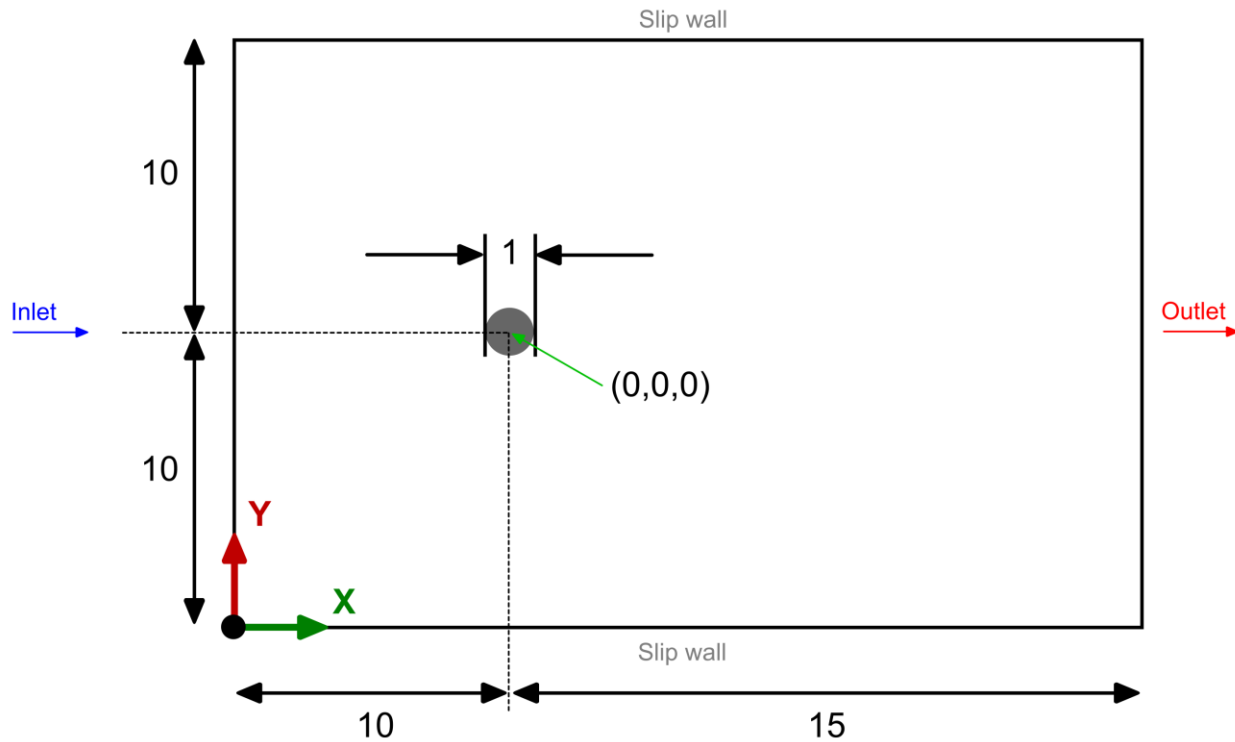


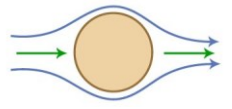
# Problem definition



All the dimensions are in meters – Figure not to scale

- Incompressible flow.
- Density:  $1 \text{ kg/m}^3$  (constant).
- Inlet velocity:  $1 \text{ m/s}$
- Adjust the viscosity to change the Reynolds number.
- Run the case for Reynolds number values ranging from 10 to 1 000 000.
  - In this range of Reynolds number values, you will encounter steady and unsteady regimes.
  - And laminar and turbulent regimes.
- Sample the lift and drag coefficients.
- Sample velocity on the wake of the cylinder.
- Choose any turbulence model with the appropriate boundary and initial conditions.

# Problem definition



**Creeping flow (no separation)**  
Steady flow

$Re < 5$



**A pair of stable vortices in the wake**  
Steady flow

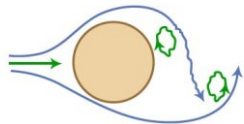
$5 < Re < 40 - 46$

Steady physics



**Laminar vortex street (Von Karman street)**  
Unsteady flow

$40 - 46 < Re < 150$



**Laminar boundary layer up to the separation point, turbulent wake**  
Unsteady flow

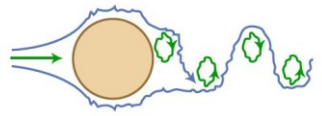
**Transition to turbulence**  
 $150 < Re < 300$   
 $300 < Re < 3 \times 10^5$

Unsteady physics



**Boundary layer transition to turbulent**  
Unsteady flow

$3 \times 10^5 < Re < 3 \times 10^6$

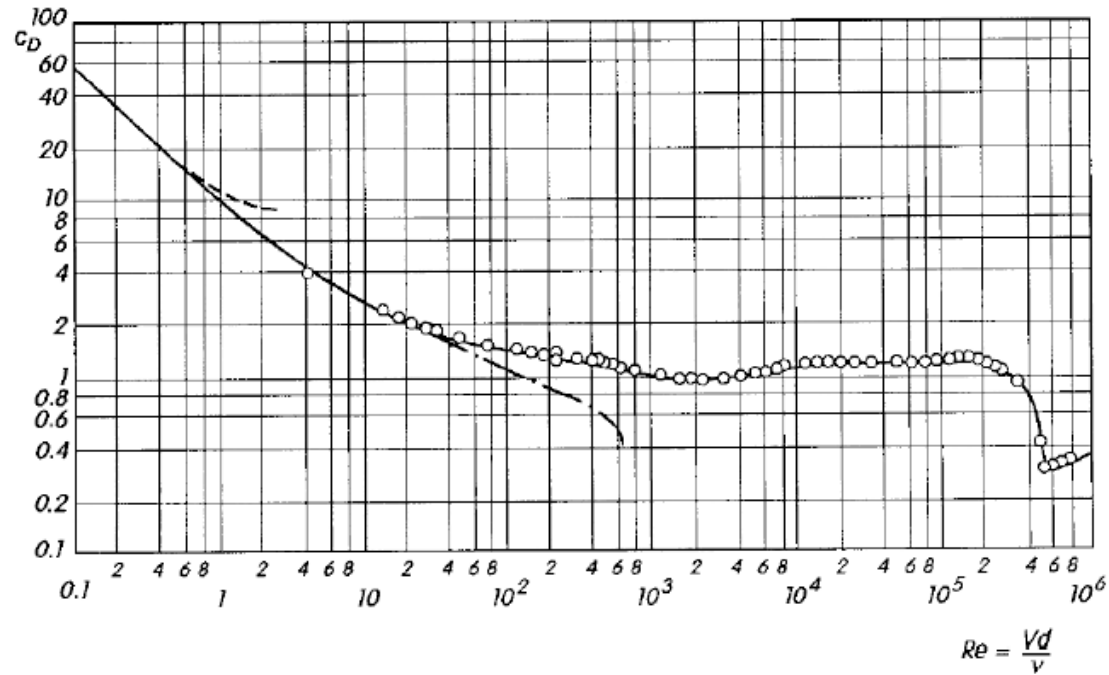


**Turbulent vortex street, but the wake is narrower than in the laminar case**  
Unsteady flow

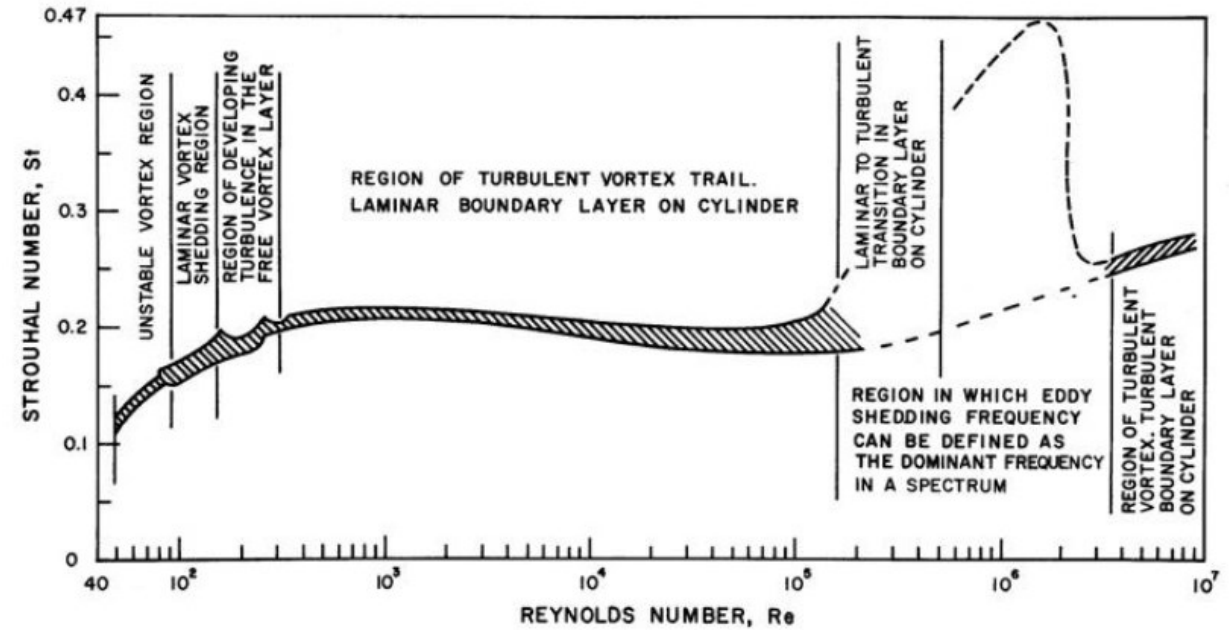
$3 \times 10^6 > Re$

# Problem definition

This case has plenty of numerical and experimental data for validation



Drag coefficient versus Reynolds Number for Circular Cylinders [1]



Strouhal Number versus Reynolds Number for Circular Cylinders [2]

[1] H. Schlichting, K. Gersten. Boundary Layer Theory. Springer, 2017.  
[2] R. Blevins. Flow-Induced Vibration. Krieger Publishing Company, 2001

# Problem definition

## Some experimental <sup>(E)</sup> and numerical <sup>(N)</sup> results of the flow past a circular cylinder at various Reynolds numbers

Reference	$c_d - Re = 20$	$L_{rb} - Re = 20$	$c_d - Re = 40$	$L_{rb} - Re = 40$
[1] Tritton <sup>(E)</sup>	2.22	–	1.48	–
[2] Cuntanceau and Bouard <sup>(E)</sup>	–	0.73	–	1.89
[3] Russel and Wang <sup>(N)</sup>	2.13	0.94	1.60	2.29
[4] Calhoun and Wang <sup>(N)</sup>	2.19	0.91	1.62	2.18
[5] Ye et al. <sup>(N)</sup>	2.03	0.92	1.52	2.27
[6] Fornberg <sup>(N)</sup>	2.00	0.92	1.50	2.24
[7] Guerrero <sup>(N)</sup>	2.20	0.92	1.62	2.21

$L_{rb}$  = length of recirculation bubble,  $c_d$  = drag coefficient,  $Re$  = Reynolds number,

[1] D. Tritton. Experiments on the flow past a circular cylinder at low Reynolds numbers. *Journal of Fluid Mechanics*, 6:547-567, 1959.

[2] M. Cuntanceau and R. Bouard. Experimental determination of the main features of the viscous flow in the wake of a circular cylinder in uniform translation. Part 1. Steady flow. *Journal of Fluid Mechanics*, 79:257-272, 1973.

[3] D. Russel and Z. Wang. A cartesian grid method for modeling multiple moving objects in 2D incompressible viscous flow. *Journal of Computational Physics*, 191:177-205, 2003.

[4] D. Calhoun and Z. Wang. A cartesian grid method for solving the two-dimensional streamfunction-vorticity equations in irregular regions. *Journal of Computational Physics*. 176:231-275, 2002.

[5] T. Ye, R. Mittal, H. Udaykumar, and W. Shyy. An accurate cartesian grid method for viscous incompressible flows with complex immersed boundaries. *Journal of Computational Physics*, 156:209-240, 1999.

[6] B. Fornberg. A numerical study of steady viscous flow past a circular cylinder. *Journal of Fluid Mechanics*, 98:819-855, 1980.

[7] J. Guerrero. Numerical simulation of the unsteady aerodynamics of flapping flight. PhD Thesis, University of Genoa, 2009.

# Problem definition

## Some experimental <sup>(E)</sup> and numerical <sup>(N)</sup> results of the flow past a circular cylinder at various Reynolds numbers

Reference	$c_d - Re = 100$	$c_l - Re = 100$	$c_d - Re = 200$	$c_l - Re = 200$
[1] Russel and Wang <sup>(N)</sup>	$1.38 \pm 0.007$	$\pm 0.322$	$1.29 \pm 0.022$	$\pm 0.50$
[2] Calhoun and Wang <sup>(N)</sup>	$1.35 \pm 0.014$	$\pm 0.30$	$1.17 \pm 0.058$	$\pm 0.67$
[3] Braza et al. <sup>(N)</sup>	$1.386 \pm 0.015$	$\pm 0.25$	$1.40 \pm 0.05$	$\pm 0.75$
[4] Choi et al. <sup>(N)</sup>	$1.34 \pm 0.011$	$\pm 0.315$	$1.36 \pm 0.048$	$\pm 0.64$
[5] Liu et al. <sup>(N)</sup>	$1.35 \pm 0.012$	$\pm 0.339$	$1.31 \pm 0.049$	$\pm 0.69$
[6] Guerrero <sup>(N)</sup>	$1.38 \pm 0.012$	$\pm 0.333$	$1.408 \pm 0.048$	$\pm 0.725$

$c_l$  = lift coefficient,  $c_d$  = drag coefficient,  $Re$  = Reynolds number

[1] D. Russel and Z. Wang. A cartesian grid method for modeling multiple moving objects in 2D incompressible viscous flow. *Journal of Computational Physics*, 191:177-205, 2003.

[2] D. Calhoun and Z. Wang. A cartesian grid method for solving the two-dimensional streamfunction-vorticity equations in irregular regions. *Journal of Computational Physics*. 176:231-275, 2002.

[3] M. Braza, P. Chassaing, and H. Hinh. Numerical study and physical analysis of the pressure and velocity fields in the near wake of a circular cylinder. *Journal of Fluid Mechanics*, 165:79-130, 1986.

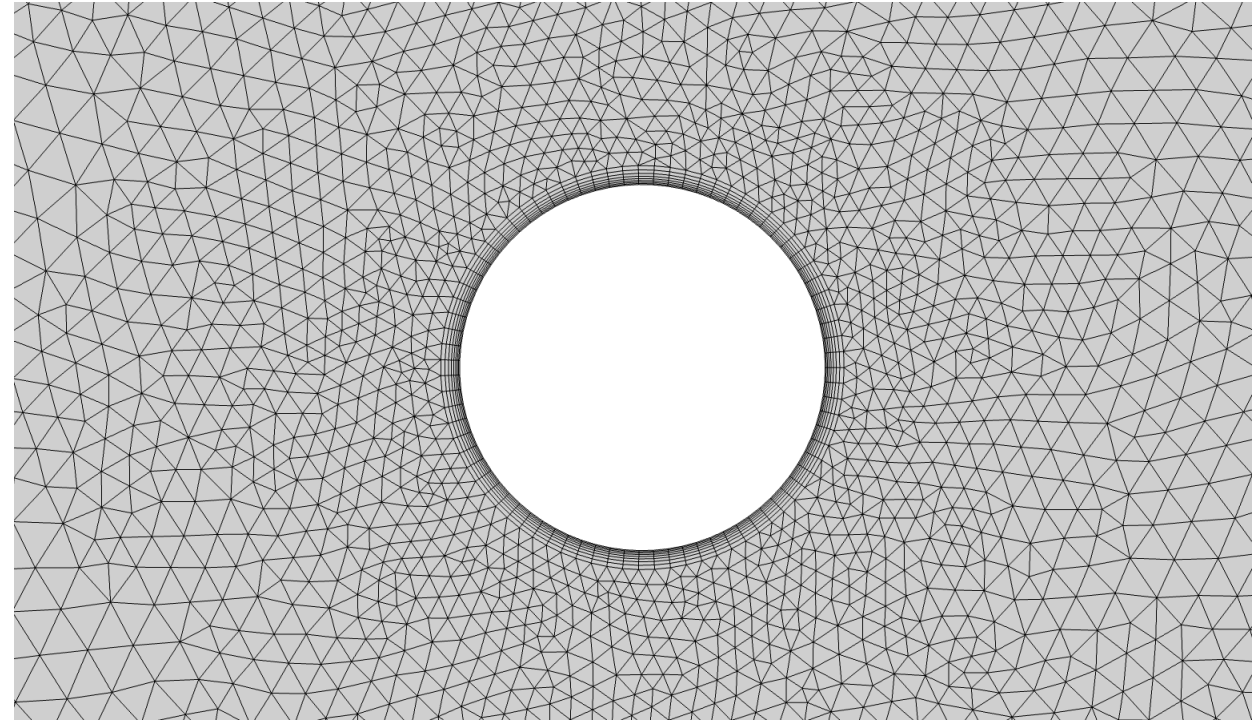
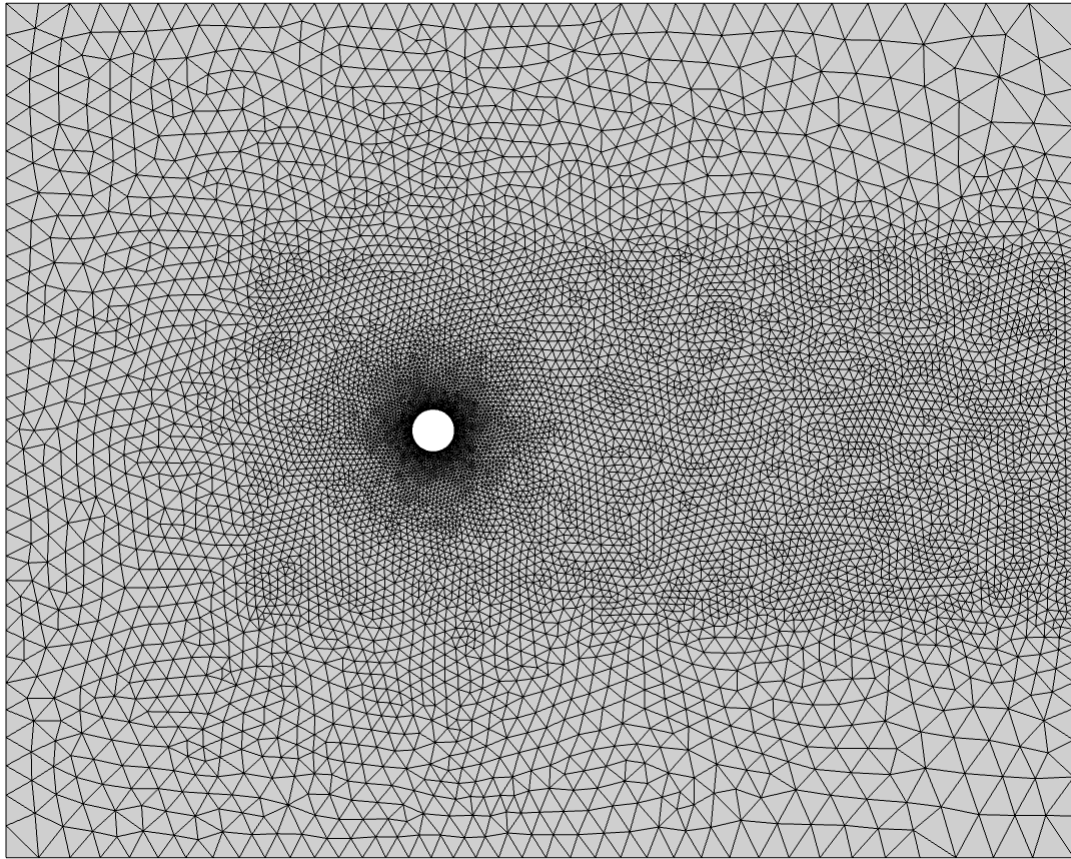
[4] J. Choi, R. Oberoi, J. Edwards, and J. Rosati. An immersed boundary method for complex incompressible flows. *Journal of Computational Physics*, 224:757-784, 2007.

[5] C. Liu, X. Zheng, and C. Sung. Preconditioned multigrid methods for unsteady incompressible flows. *Journal of Computational Physics*, 139:33-57, 1998.

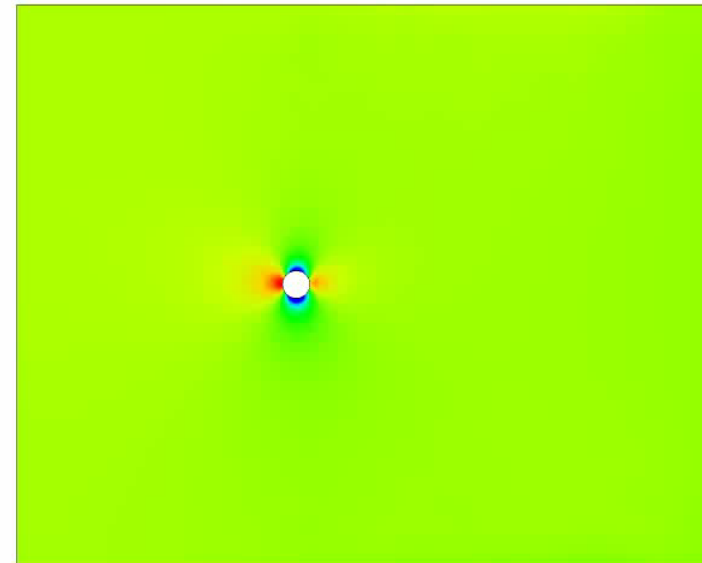
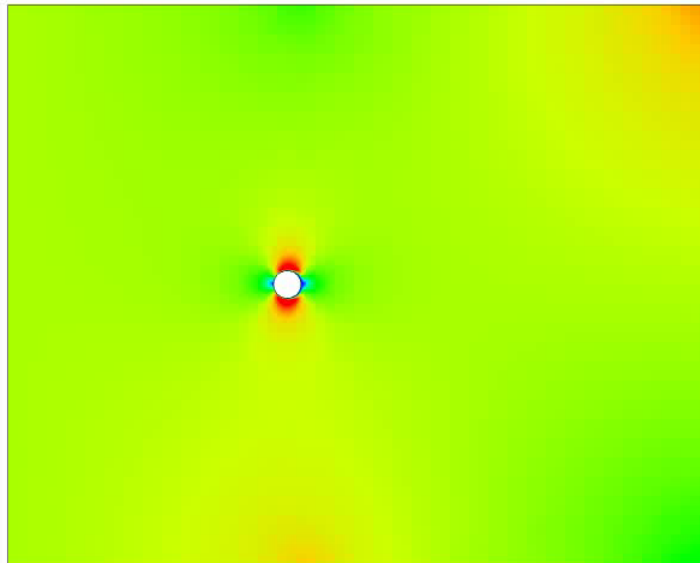
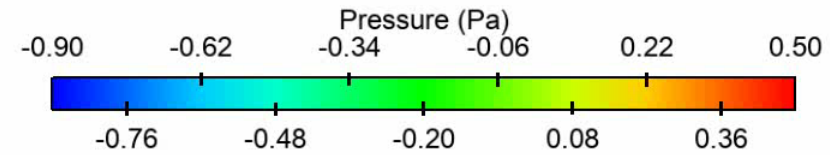
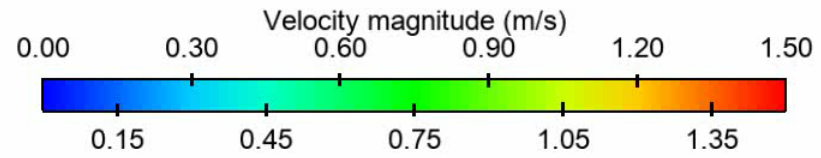
[6] J. Guerrero. Numerical Simulation of the unsteady aerodynamics of flapping flight. PhD Thesis, University of Genoa, 2009.

# Domain mesh – 2D mesh and inflation layer around the cylinder

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# Qualitative results – Contour plots



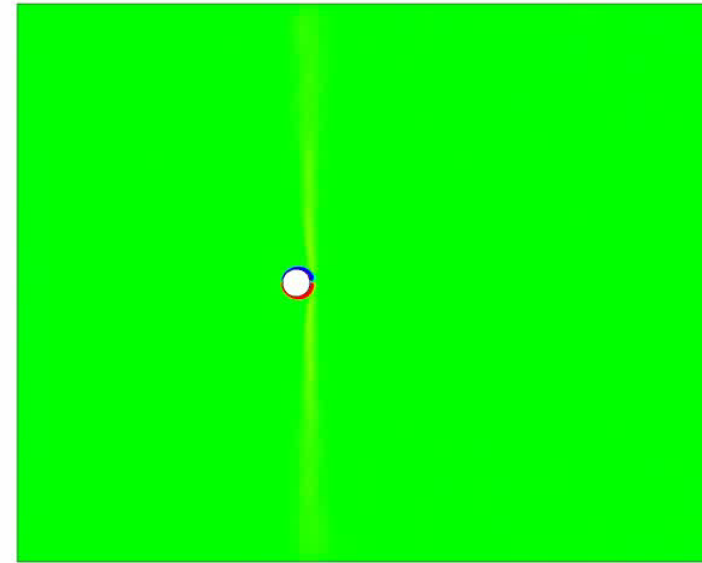
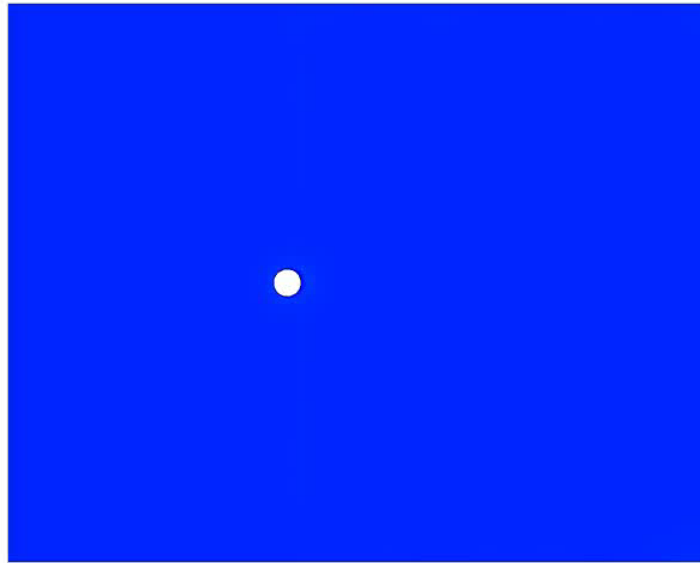
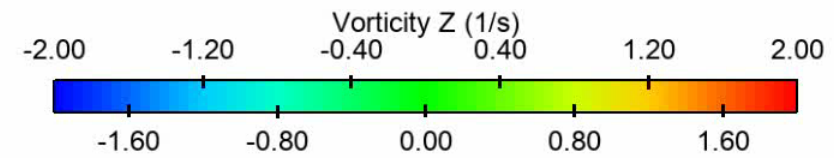
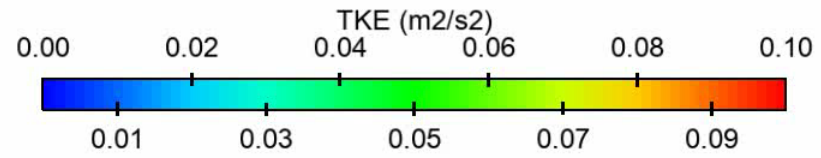
Time = 0.50 s

In all cases the Reynolds number is in the unsteady and turbulent regime.

**Left:** velocity magnitude contours – **Right:** relative pressure contours.

[www.wolfdynamics.com/training/turbulence/movies1/mov1.avi](http://www.wolfdynamics.com/training/turbulence/movies1/mov1.avi)

# Qualitative results – Contour plots



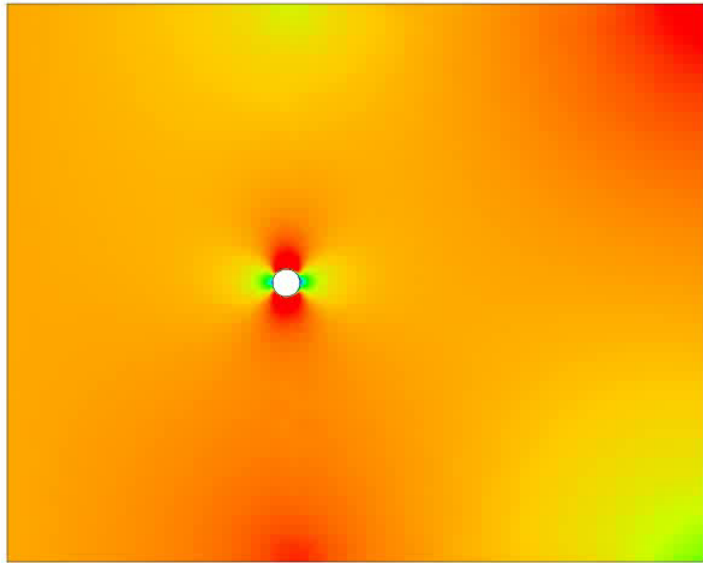
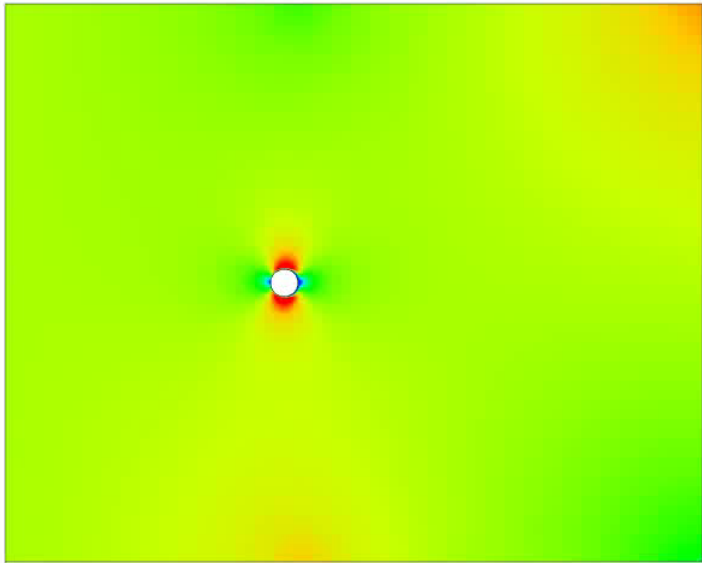
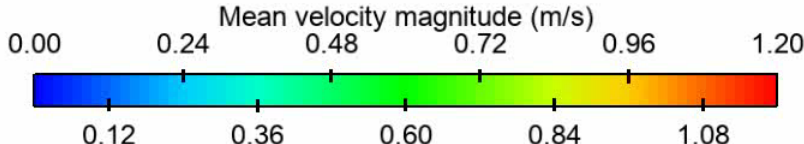
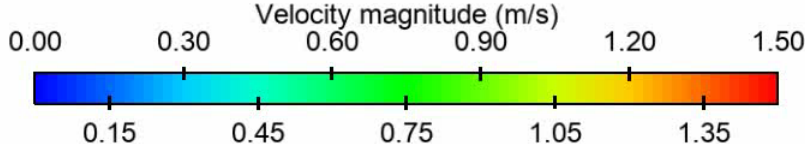
Time = 0.50 s

**Left:** turbulent kinetic energy contours – **Right:** vorticity Z contours (component normal to the screen).

[www.wolfdynamics.com/training/turbulence/movies1/mov2.avi](http://www.wolfdynamics.com/training/turbulence/movies1/mov2.avi)



# Qualitative results – Contour plots

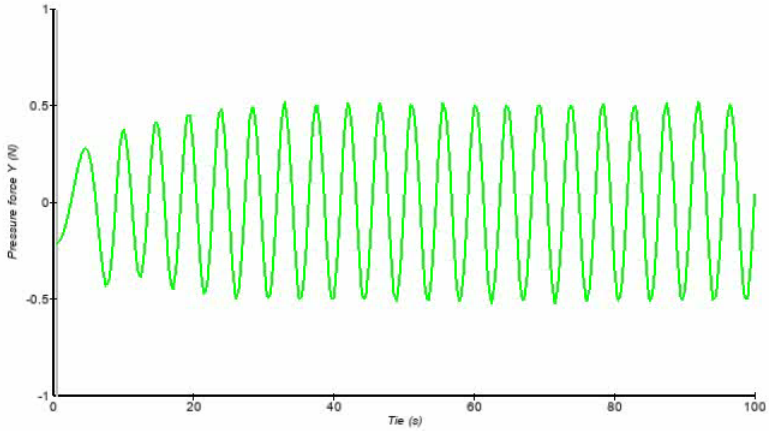
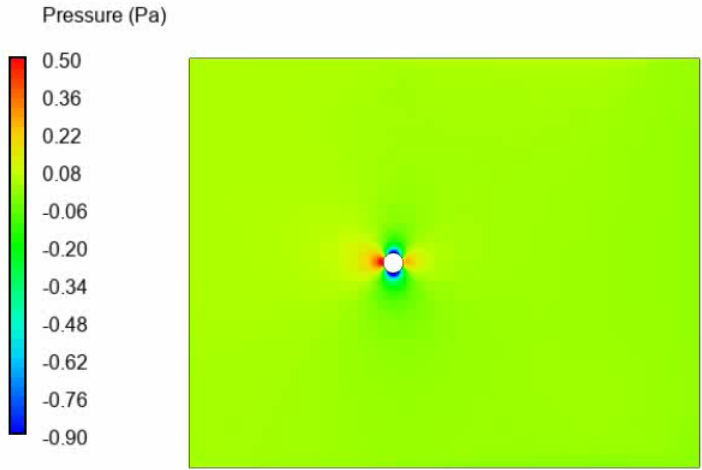
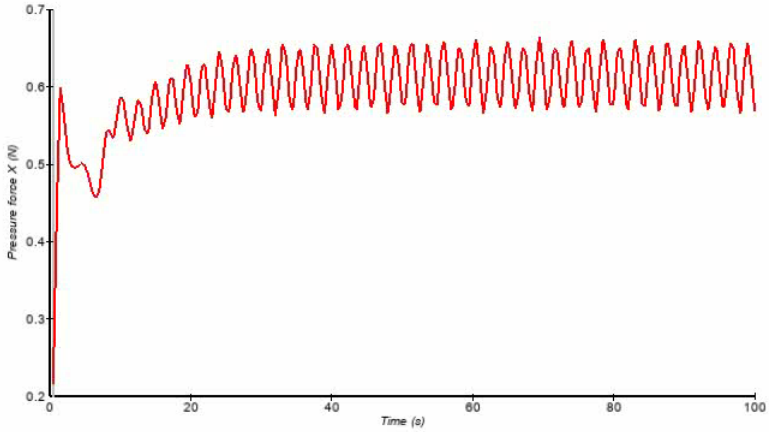
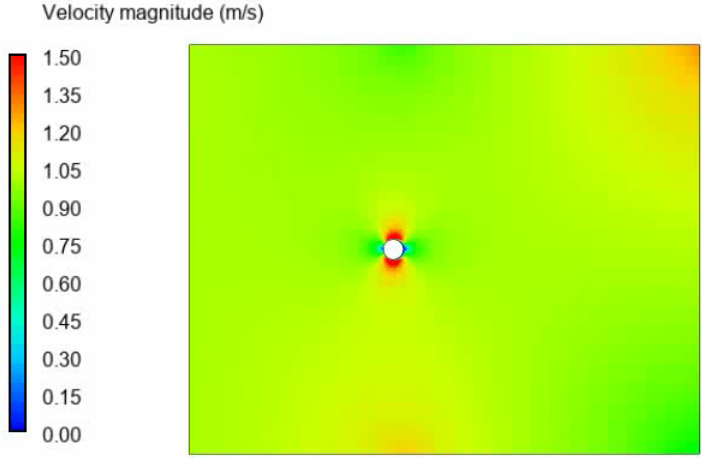


Time = 0.50 s

**Left:** velocity magnitude contours – **Right:** mean velocity magnitude contours

[www.wolfdynamics.com/training/turbulence/movies1/mov3.avi](http://www.wolfdynamics.com/training/turbulence/movies1/mov3.avi)

# Qualitative results – Contour plots

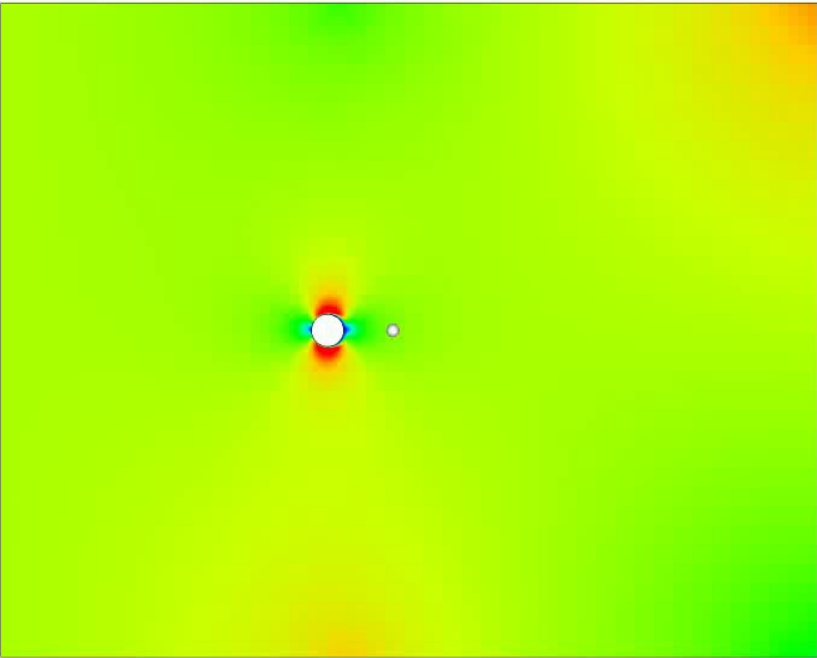
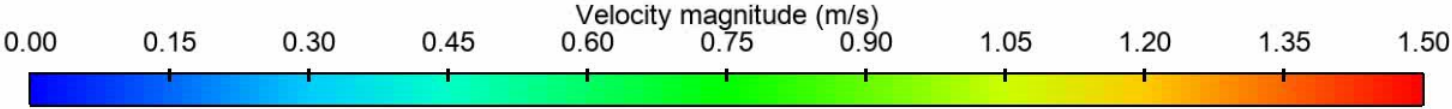


Time = 0.50 s

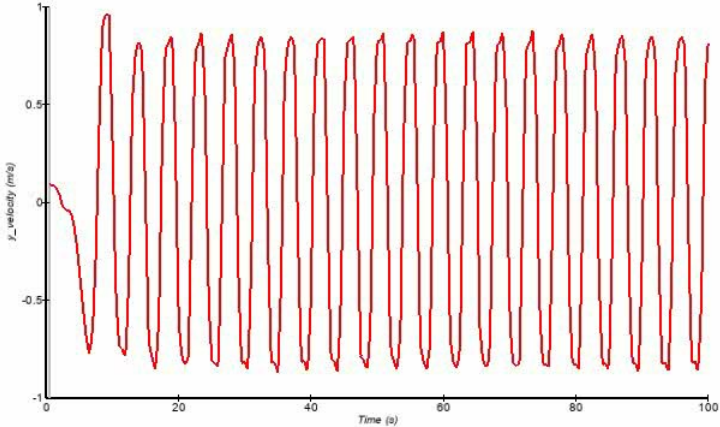
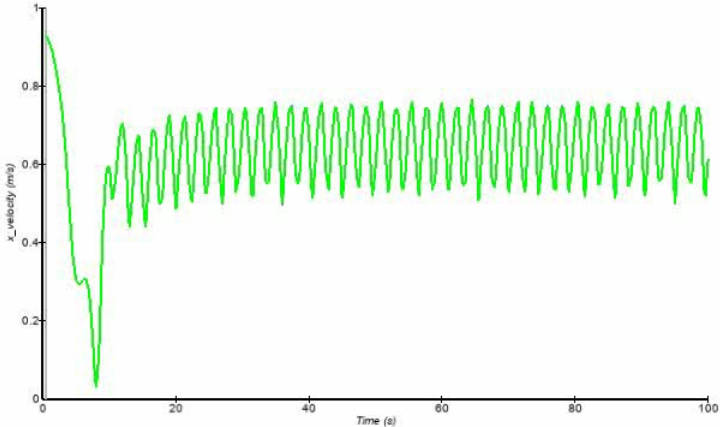
**Left:** velocity magnitude and pressure contours – **Right:** Components X and Y of pressure force over the cylinder

[www.wolfdynamics.com/training/turbulence/movies1/mov4.avi](http://www.wolfdynamics.com/training/turbulence/movies1/mov4.avi)

# Qualitative results – Contour plots



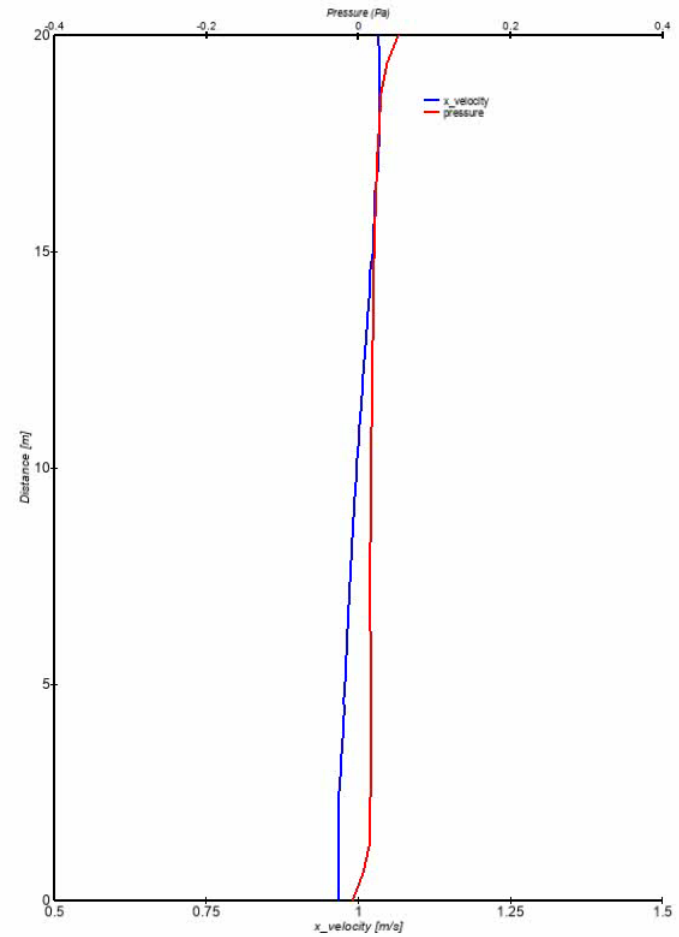
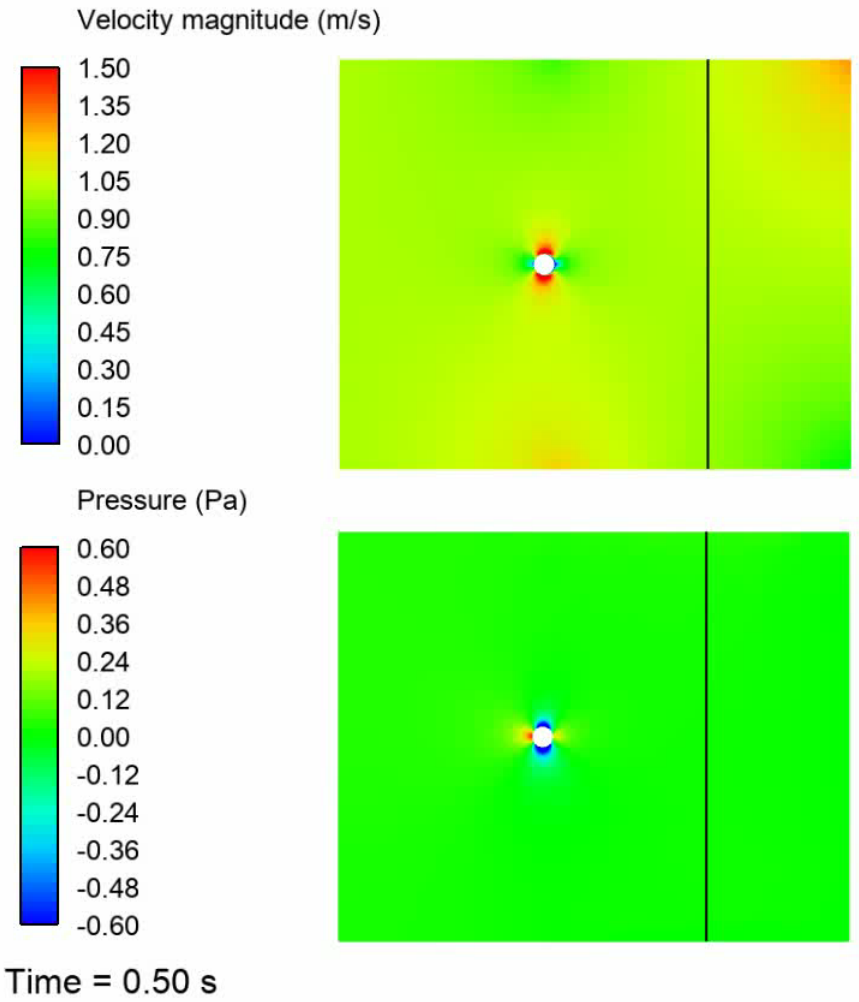
Time = 0.50 s



**Left:** velocity magnitude contours – **Right:** velocity sampled at a location in the wake of the cylinder

[www.wolfdynamics.com/training/turbulence/movies1/mov5.avi](http://www.wolfdynamics.com/training/turbulence/movies1/mov5.avi)

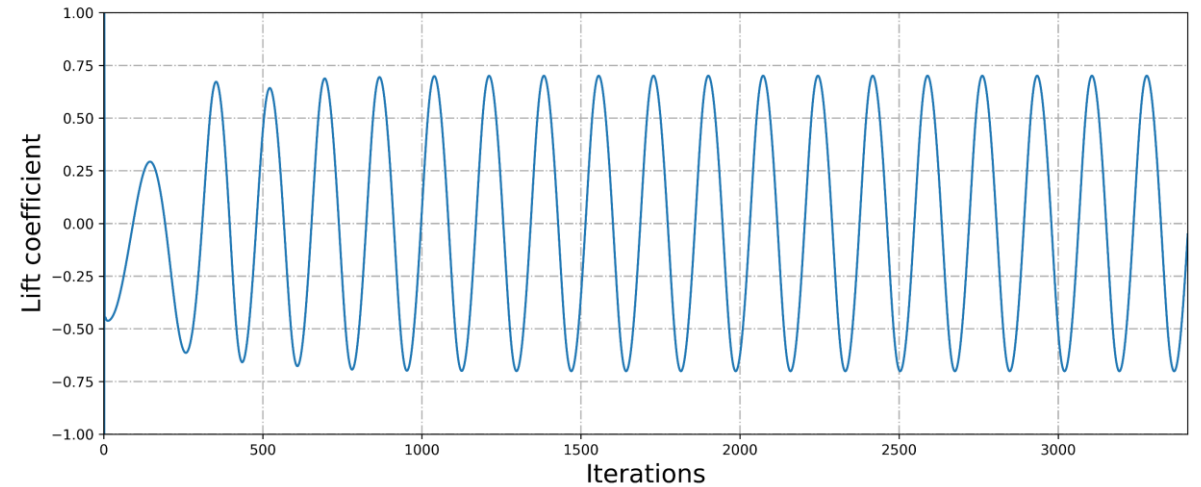
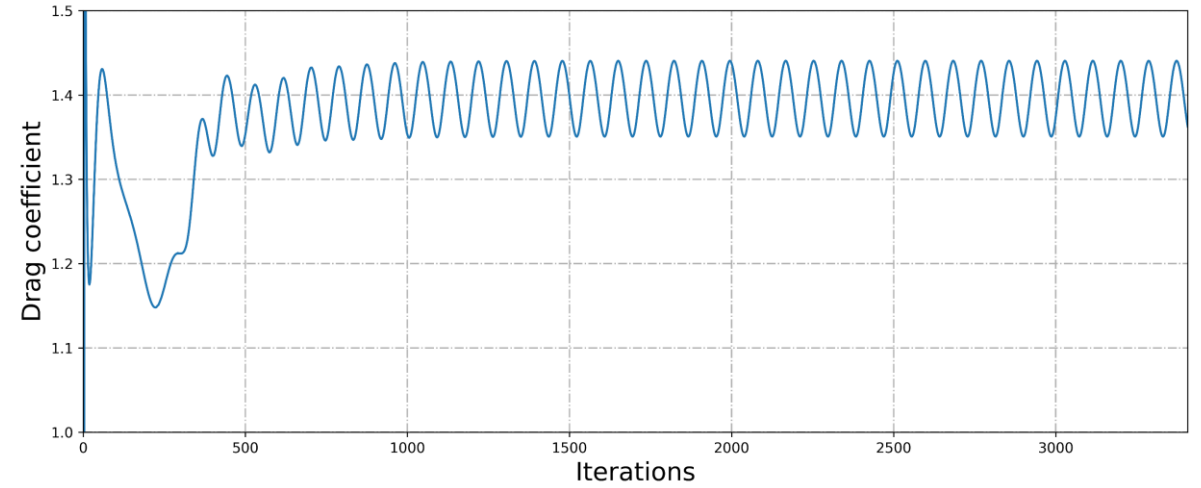
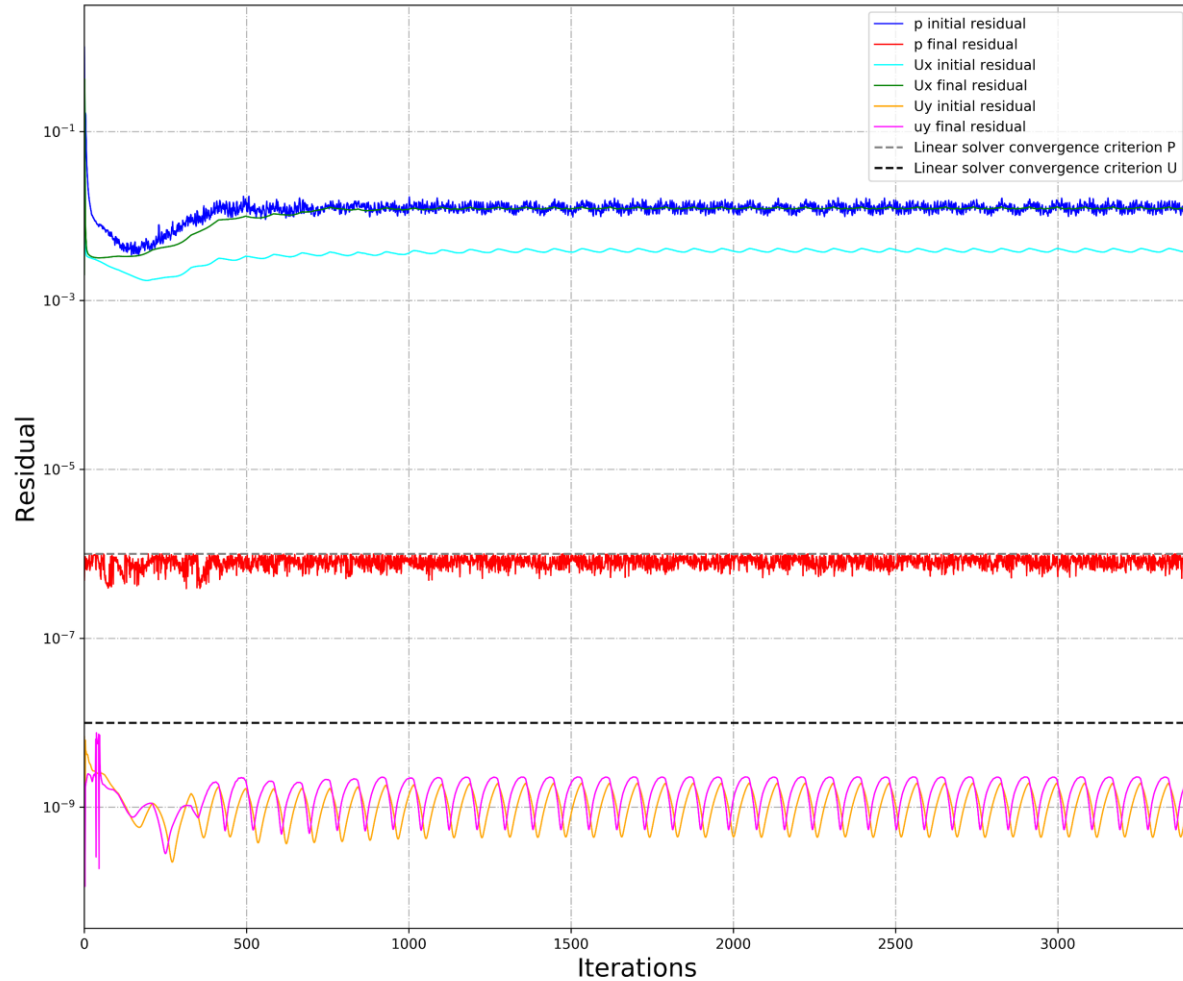
# Qualitative results – Contour plots



**Left:** velocity magnitude contours – **Right:** velocity and pressure sample in a vertical line in the wake of the cylinder

[www.wolfdynamics.com/training/turbulence/movies1/mov7.avi](http://www.wolfdynamics.com/training/turbulence/movies1/mov7.avi)

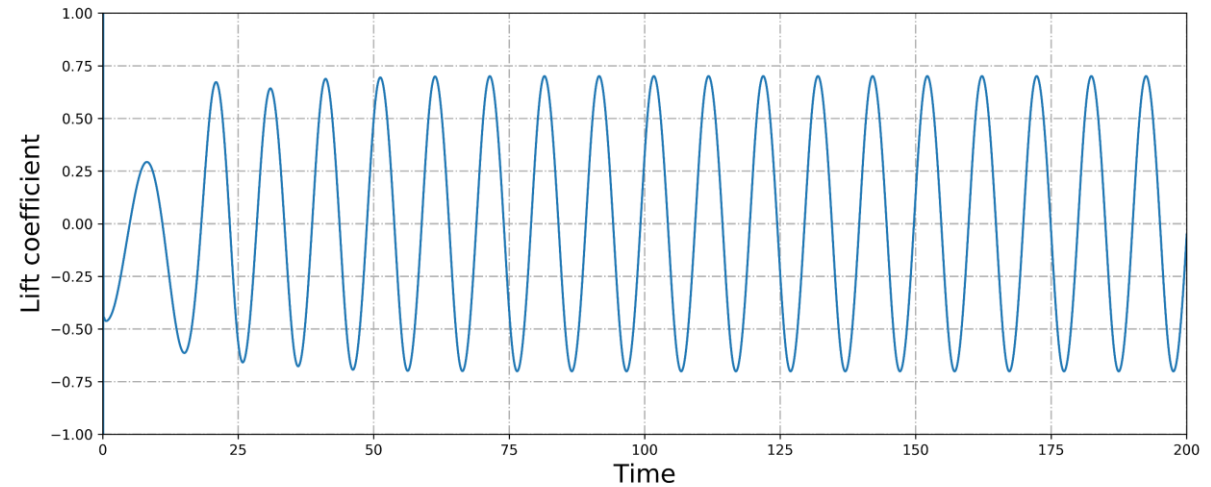
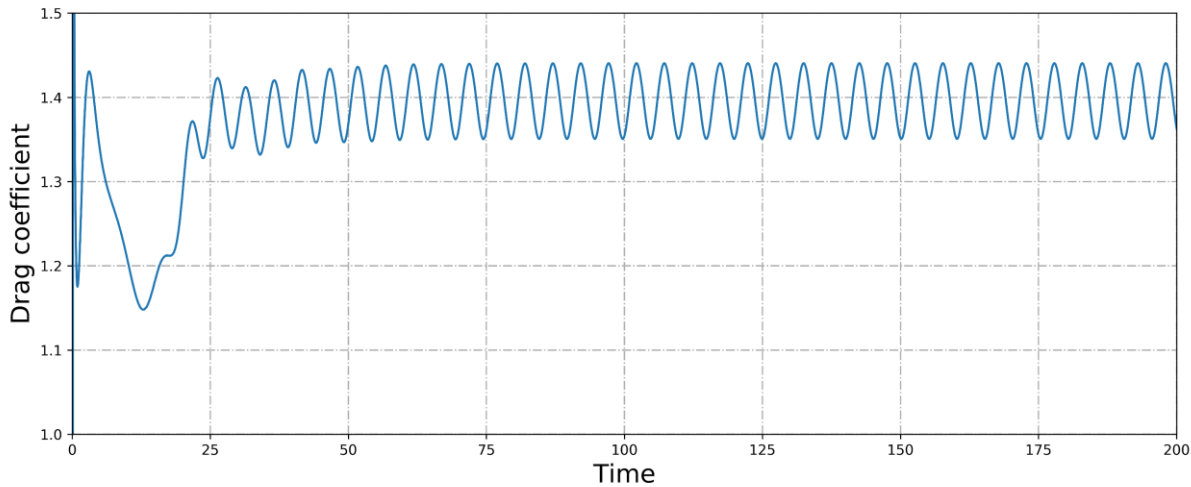
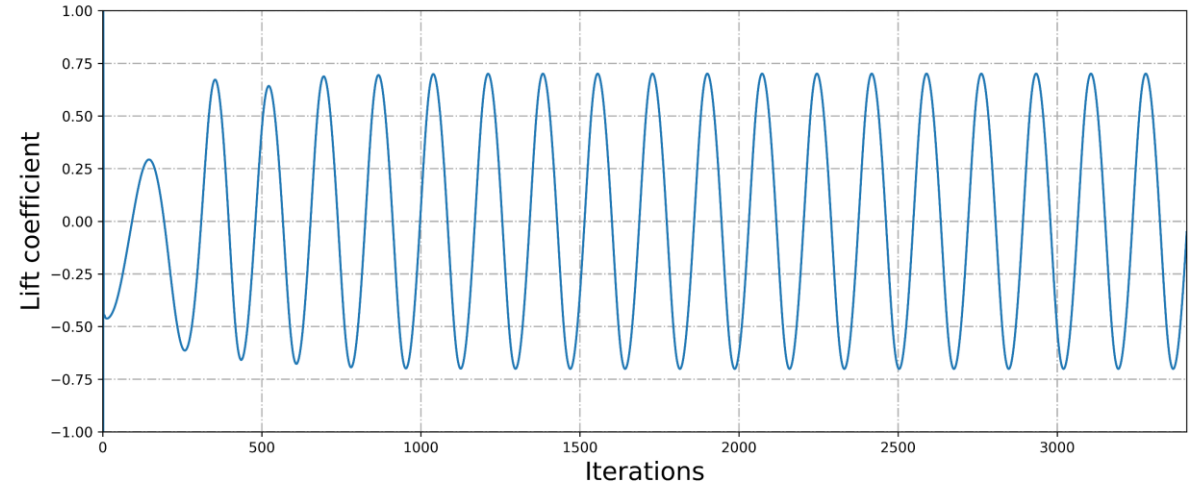
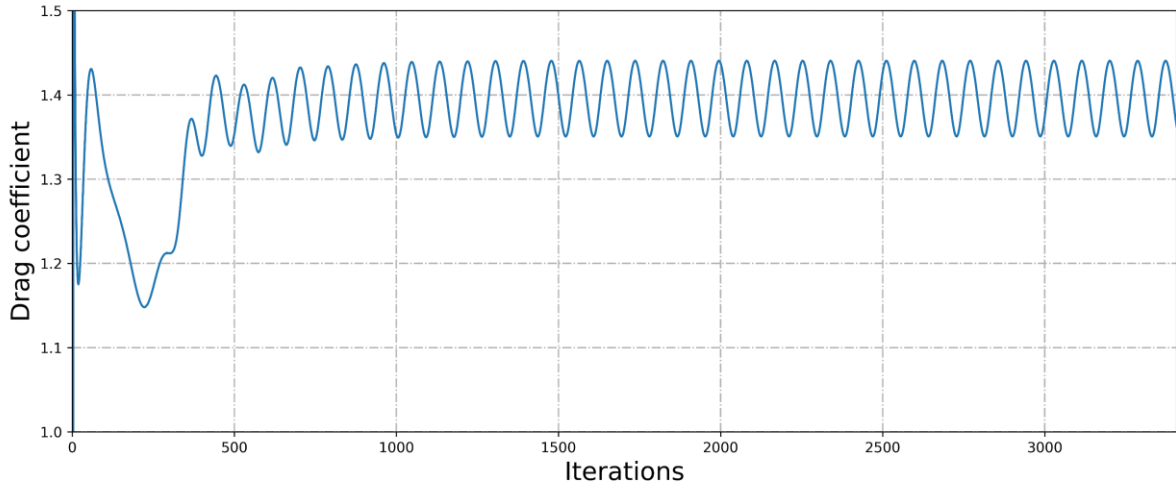
# Quantitative results – Residuals – Lift and drag time series



**Unsteady solver – Comparison of iterative convergence and iterative QoI**

Lift coefficient and drag coefficient signals

# Quantitative results – Lift and drag time series



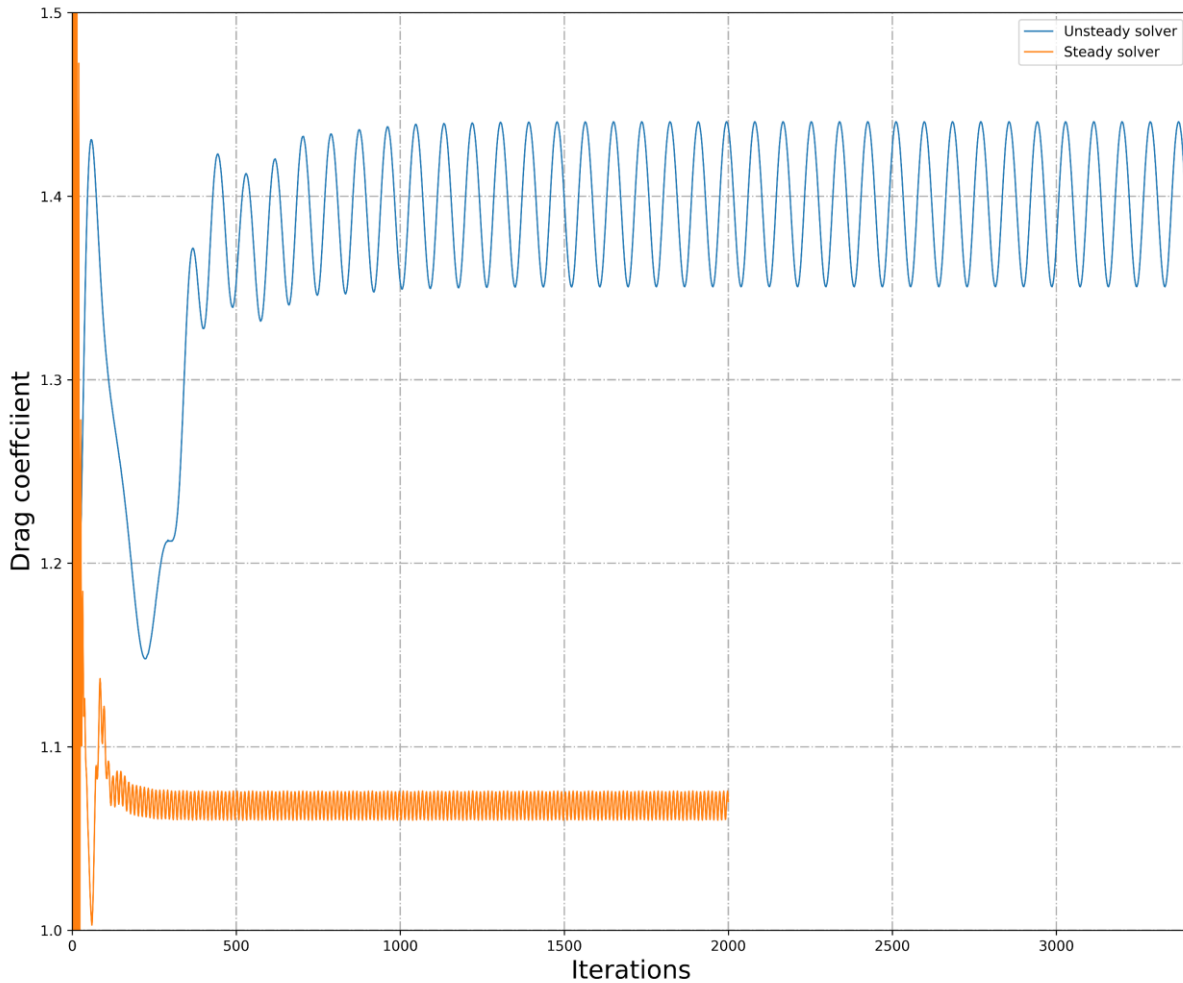
## Unsteady solver QoI – Drag coefficient

Top image: in function of iterations  
Bottom image: in function of time

## Unsteady solver QoI – Lift coefficient

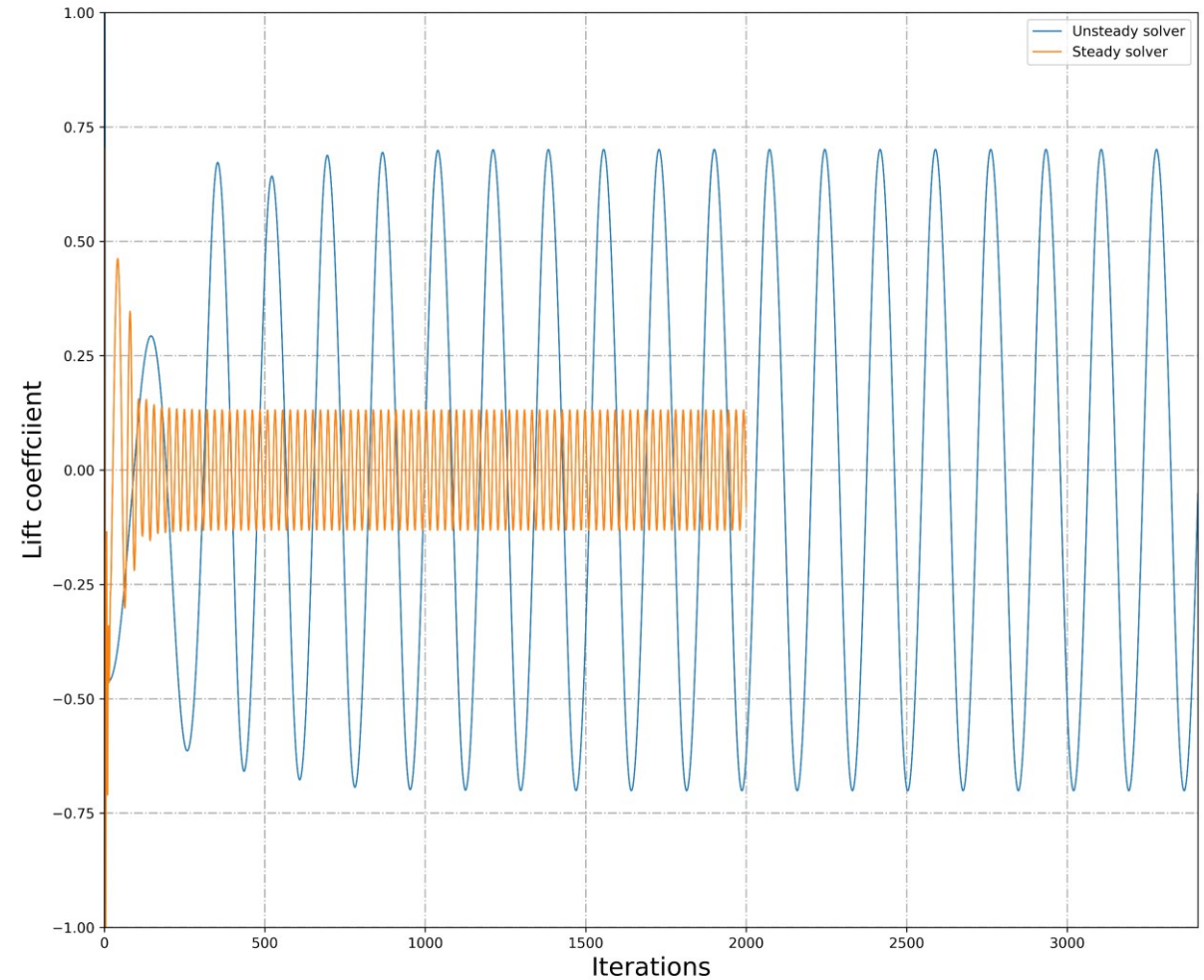
Top image: in function of iterations  
Bottom image: in function of time

# Quantitative results – Lift and drag time series



**Unsteady solver and steady solver comparison**

Comparison of the drag coefficient outcome

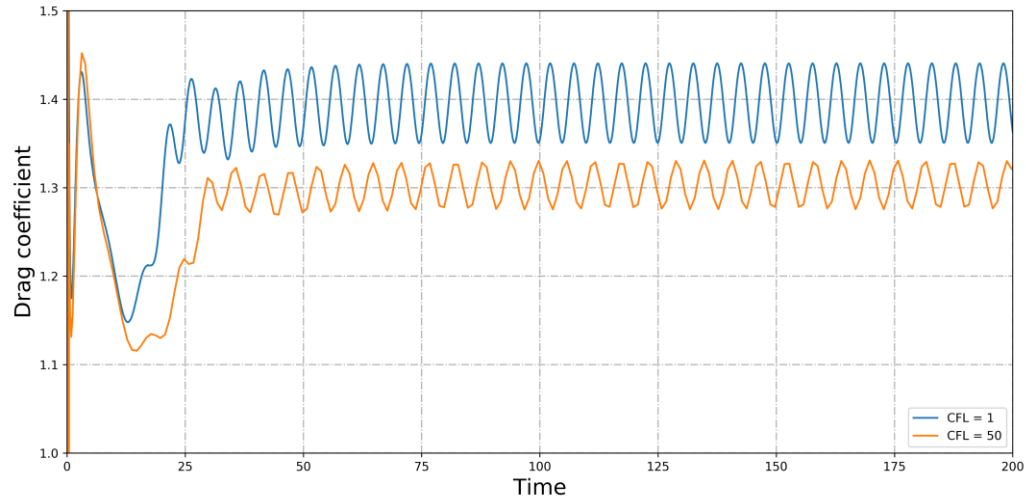


**Unsteady solver and steady solver comparison**

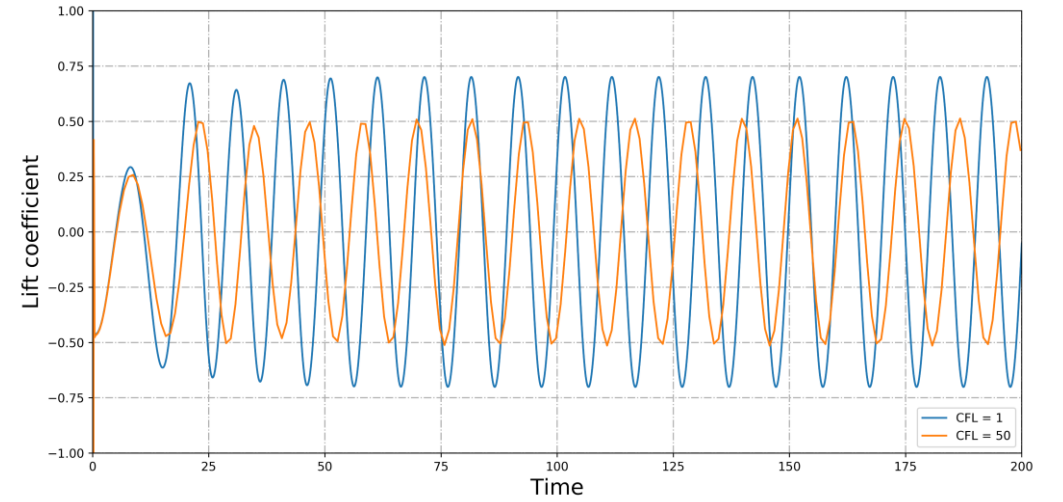
Comparison of the lift coefficient outcome



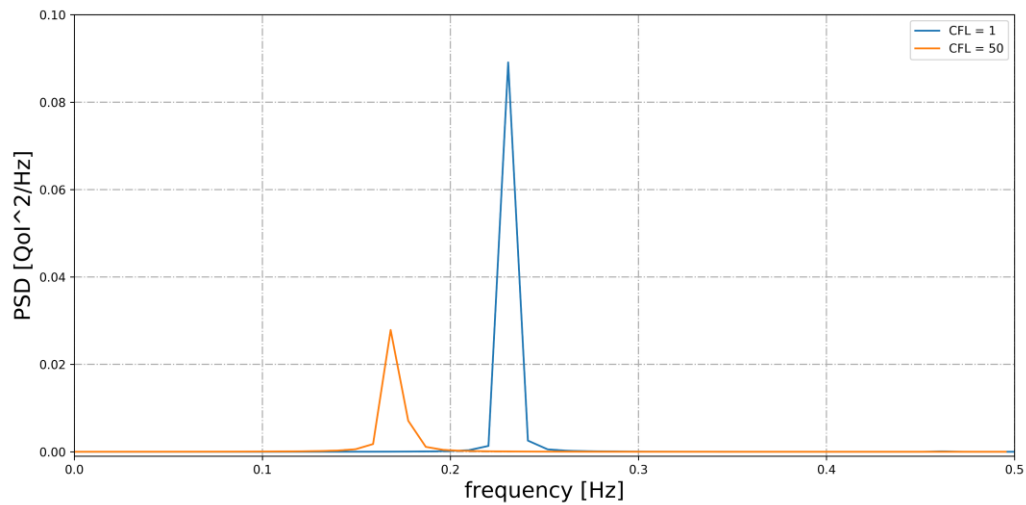
# Qualitative results – Contour plots



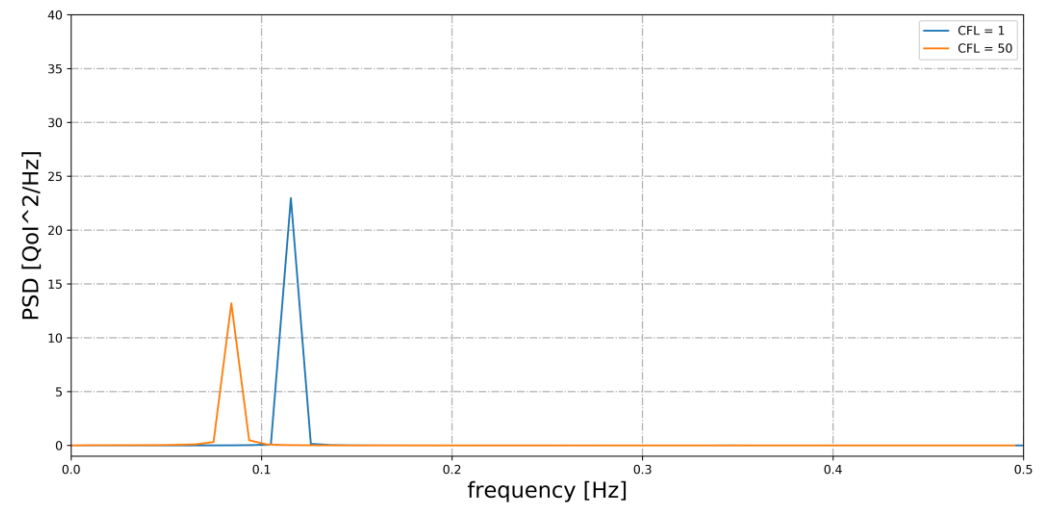
Drag coefficient time series  
CFL = 1 and CFL = 50



Lift coefficient time series  
CFL = 1 and CFL = 50



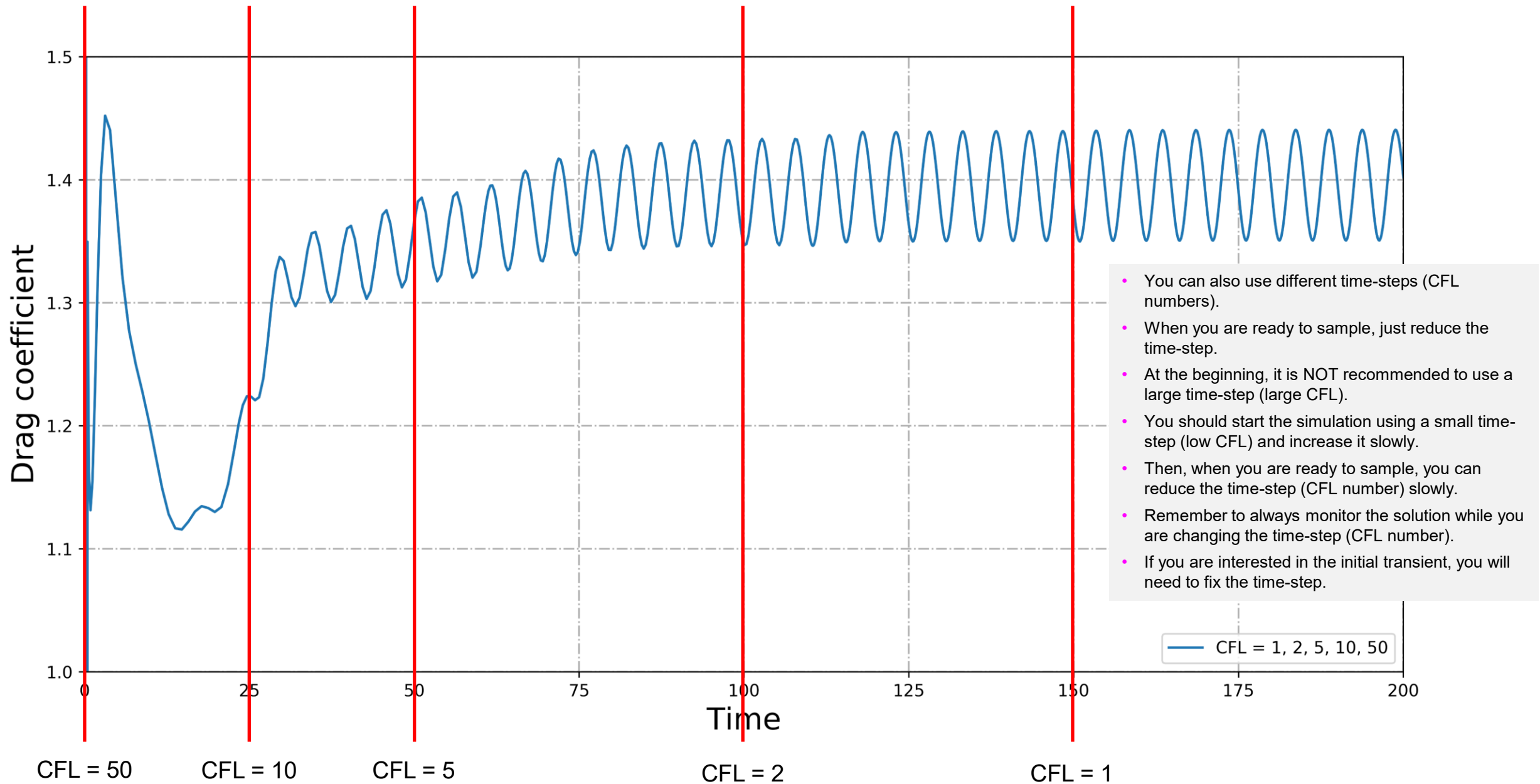
Periodogram of drag coefficient time series  
CFL = 1 and CFL = 50



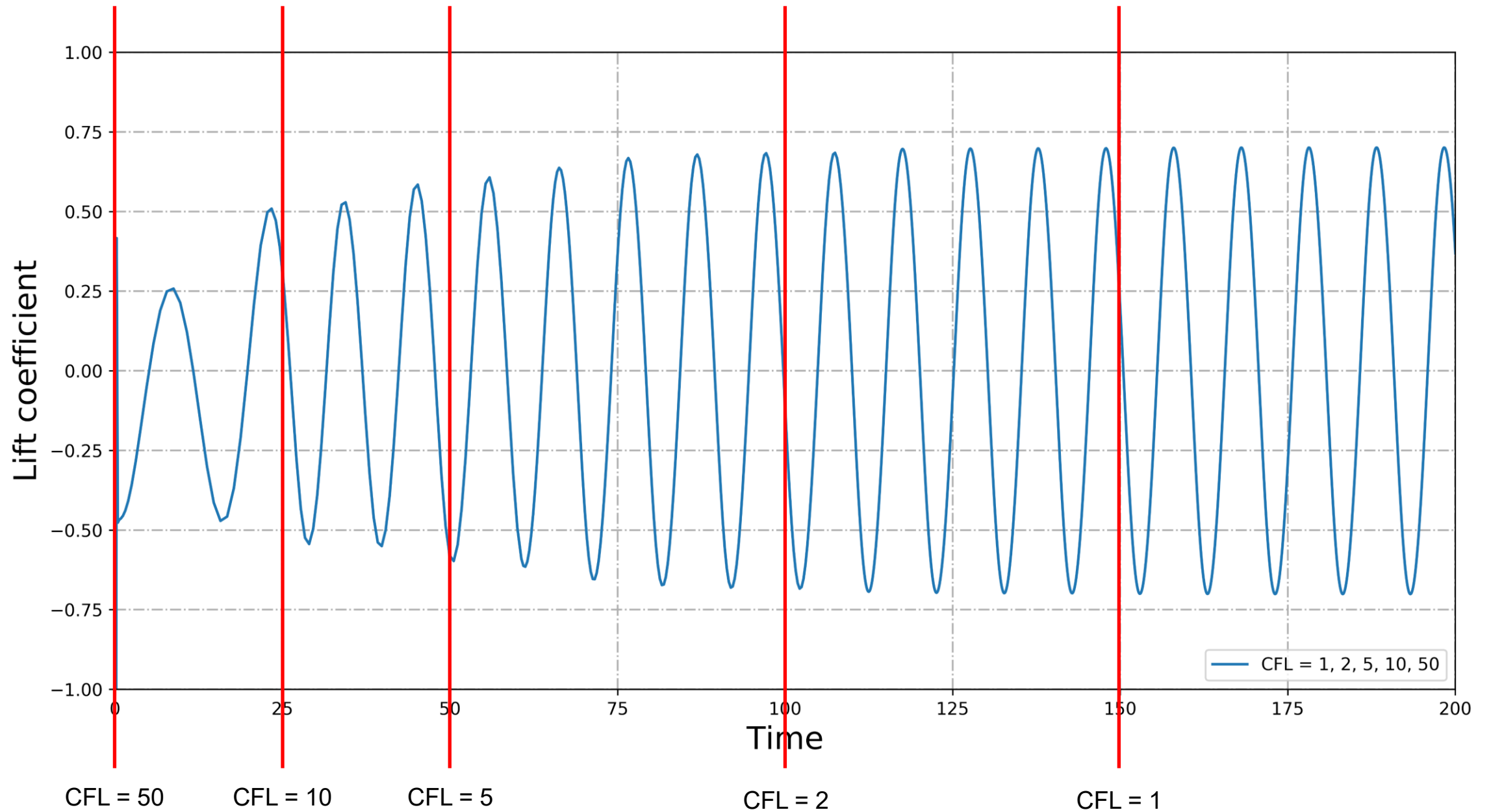
Periodogram of lift coefficient time series  
CFL = 1 and CFL = 50



# Qualitative results – Contour plots



# Qualitative results – Contour plots



# On the CFL number

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- First of all, what is the **CFL number** or **Courant number**?
- In one dimension, the **CFL number** is defined as,

$$CFL = \frac{u \Delta t}{\Delta x}$$


- The **CFL number** is a measure of how much information ( $u$ ) traverses a computational grid cell ( $\Delta x$ ) in a given time-step ( $\Delta t$ ).

# On the CFL number

---

- Similarly, we can define the **CFL number condition**, which is related to the **CFL number**.
- For the **N** dimensional case, the **CFL number condition** becomes,

$$CFL = \Delta t \sum_{i=1}^n \frac{u_i}{\Delta x_i} \leq CFL_{max}$$



Maximum CFL number allowed  
by the numerical method

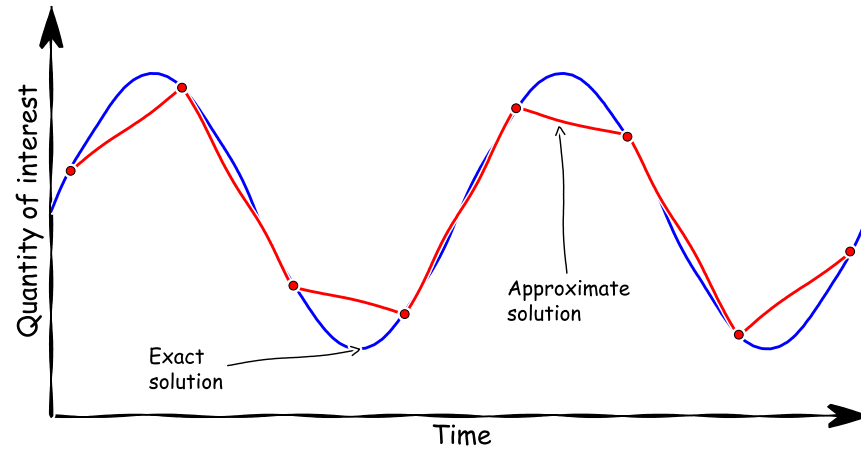
# On the CFL number

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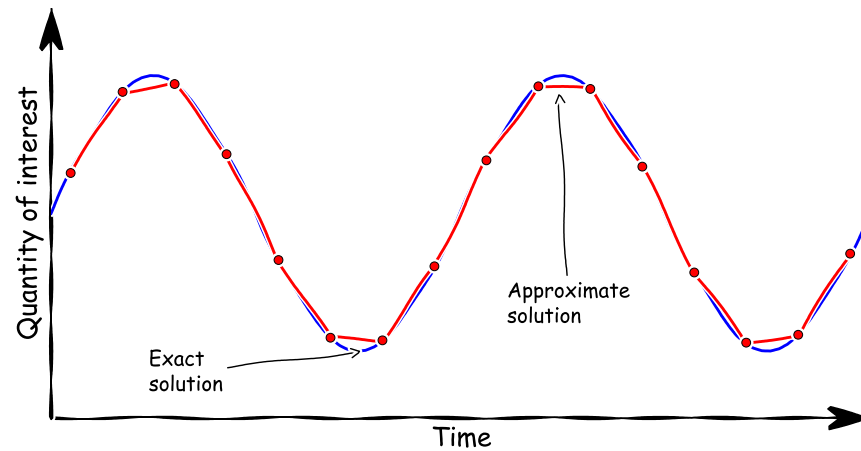
- The **CFL number** is a necessary condition to guarantee the stability of the numerical scheme.
- But not all numerical schemes have the same stability requirements or  $CFL_{\max}$  requirement (maximum allowable **CFL number**).
- We are going to use implicit numerical methods which are **unconditionally stable**.
- In other words, they are not constrained to the **CFL number condition** or maximum allowable **CFL number**.
- However, the fact that we are going to use a numerical method that is unconditionally stable, does not mean that we can use a time step of any size.
- The time-step must be chosen in such a way that it resolves the time-dependent features, and it maintains the solver stability.

# On the CFL number

- When running unsteady simulations, the time-step must be chosen in such a way that it resolves the time-dependent features and maintains solver stability.



When you use large time steps you do not resolve well the physics



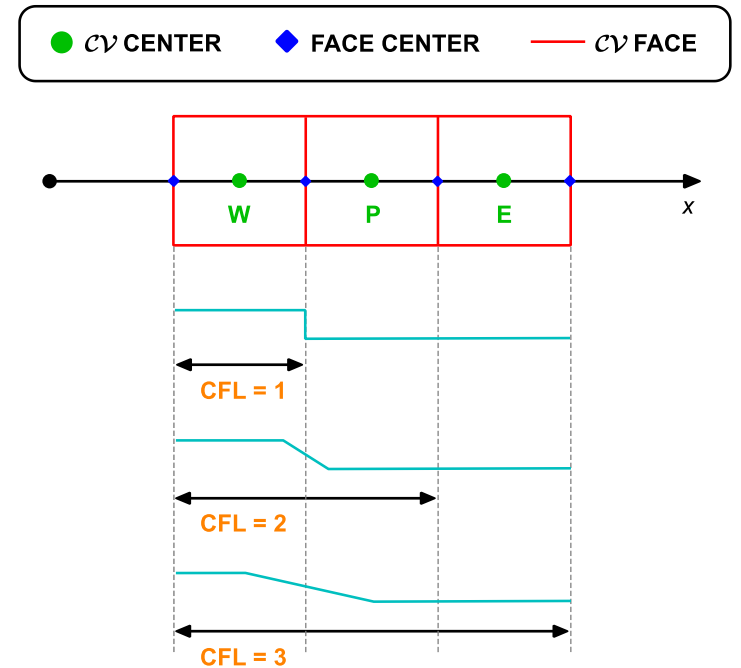
By using a smaller time step you resolve better the physics, and you gain stability

# On the CFL number

- I like to see the **CFL number** as follows,

$$CFL = \frac{u \Delta t}{\Delta x} = \frac{u}{\Delta x / \Delta t} = \frac{\text{speed of the PDE}}{\text{speed of the mesh}}$$

- It is an indication of the amount of information that propagates through one cell (or many cells), in one time-step.



- The **CFL condition** is a necessary condition for stability (and hence convergence).
- But it is not always sufficient to guarantee stability.
- Other properties of the discretization schemes that you should observe are: conservationness, boundedness, transportiveness, and accuracy.
- The **CFL number** is not a magical number.