# Optimal control of complex flows 

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The work presented has been carried out by:

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- Tom Bewley UCSD, USA;
(1) Introduction
(2) MCE: Minimal Control Energy
(3) ADA: Adjoint of the Direct-Adjoint

4. Applications
(5) Numerics

## Introduction

## Definitions

Complex flows
here problems with large number of degrees of freedom
Optimal control
The linearized system

$$
\frac{\partial \mathbf{x}}{\partial t}=A \mathbf{x}+B \mathbf{u} \quad \text { on } \quad 0<t<T, \quad \text { with } \quad \mathbf{x}=\mathbf{x}_{0} \quad \text { at } \quad t=0
$$

- $\mathbf{x}$ has dimension $n$ and $\mathbf{u}$ dimension $m$
- here $n \gg m$
- find $\mathbf{u}$ that minimizes a quadratic cost function $J$
- consider: full state information, no estimation (has been done, references available)


## The linear optimal control problem

The classical full-state-information control problem is formulated as: find the control $\mathbf{u}$ that minimizes the cost function

$$
J=\frac{1}{2} \int_{0}^{T}\left[\mathbf{x}^{H} Q \mathbf{x}+I^{2} \mathbf{u}^{H} R \mathbf{u}\right] d t
$$

where $I$ is the penalty of the control, and the state x and the control u are related via the state equation

$$
\frac{\partial \mathbf{x}}{\partial t}=A \mathbf{x}+B \mathbf{u} \quad \text { on } \quad 0<t<T, \quad \text { with } \quad \mathbf{x}=\mathbf{x}_{0} \quad \text { at } \quad t=0 .
$$

The solution depends on: $x_{0}, T, Q, R$ and $/$.

## Solution approaches

- With a feedback rule $\mathbf{u}=K \mathbf{x}$, and a system which is LTI, then the feedback matrix $K$ is computed once off-line (convenient since $K$ is independent of $\mathbf{x}_{0}$ ).
- Optimal control $\mathbf{u}$ corresponding to the state at each time step is computed in real time, normally with a finite horizon (value of $T$ ) to make it tractable. (Example: Adjoint-based control optimization.)

Both approaches can be solved using the adjoint of the state equation.

## Why introduce Adjoint equations ?

Example gradient computation:

$$
J=\mathbf{w}^{H} \mathbf{x}, \quad \text { where } \quad A \mathbf{x}=\mathbf{b}, \quad \text { Ex. find } \quad \frac{\partial J}{\partial \mathbf{b}}
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Finite difference approach

$$
\left(\frac{\partial J}{\partial \mathbf{b}}\right)_{j} \approx \frac{J\left(\mathbf{b}+\epsilon \mathbf{e}_{j}\right)-J(\mathbf{b})}{\epsilon}
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then

$$
J=\mathbf{w}^{H} \mathbf{x}=\left(A^{H} \mathbf{p}\right)^{H} \mathbf{x}=\mathbf{p}^{H} A \mathbf{x}=\mathbf{p}^{H} \mathbf{b}
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$$

and

$$
\frac{\partial J}{\partial \mathbf{b}}=\mathbf{p} \quad \text { One Solution, independently of } n
$$

## Derivation of adjoint I

The adjoint variable $\mathbf{p}$ is introduced as a Lagrange multiplier. The augmented cost function is written

$$
J=\int_{0}^{T} \frac{1}{2}\left[\mathbf{x}^{H} Q \mathbf{x}+I^{2} \mathbf{u}^{H} R \mathbf{u}\right]-\mathbf{p}^{H}\left[\frac{\partial \mathbf{x}}{\partial t}-A \mathbf{x}-B \mathbf{u}\right] d t
$$

linearize + integration by parts and $\delta J=0$ gives

$$
0=\int_{0}^{T} \delta \mathbf{u}^{H}\left[B \mathbf{p}+I^{2} R \mathbf{u}\right]+\delta \mathbf{x}^{H}\left[\frac{\partial \mathbf{p}}{\partial t}+A^{H} \mathbf{p}+Q \mathbf{x}\right] d t+\left[\delta \mathbf{x}^{H} \mathbf{p}\right]_{0}^{T},
$$

## Derivation of adjoint II

The adjoint variable $\mathbf{p}$ is introduced as a Lagrange multiplier. The augmented cost function is written

$$
J=\int_{0}^{T} \frac{1}{2}\left[\mathbf{x}^{H} Q \mathbf{x}+I^{2} \mathbf{u}^{H} R \mathbf{u}\right]-\mathbf{p}^{H}\left[\frac{\partial \mathbf{x}}{\partial t}-A \mathbf{x}-B \mathbf{u}\right] d t
$$

linearize + integration by parts and $\delta J=0$

$$
0=\int_{0}^{T} \delta \mathbf{u}^{H}\left[B \mathbf{p}+I^{2} R \mathbf{u}\right]+\delta \mathbf{x}^{H}[\underbrace{\frac{\partial \mathbf{p}}{\partial t}+A^{H} \mathbf{p}+Q \mathbf{x}}_{=0}] d t+\left[\delta \mathbf{x}^{H} \mathbf{p}\right]_{0}^{T},
$$

gives adjoint equations (obs! $\delta \mathbf{x}(0)=0)$

$$
\frac{\partial \mathbf{p}}{\partial t}=-A^{H} \mathbf{p}-Q \mathbf{x}, \quad \text { with } \quad \mathbf{p}(t=T)=0
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linearize + integration by parts and $\delta J=0$

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$$

and optimality condition

$$
\mathbf{u}=-\frac{1}{l^{2}} R^{-1} B^{H} \mathbf{p} .
$$

## Optimal control using feedback

If we consider a feedback rule $\mathbf{u}=K \mathbf{x}$ then

$$
\mathbf{u}=K \mathbf{x}=-\frac{1}{l^{2}} R^{-1} B^{H} \mathbf{p} .
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## How does it work ?

Note that state is often denoted direct

## Two-point boundary value problem

Write the direct and adjoint equations on a combined matrix form

$$
\begin{gathered}
\frac{d \mathbf{z}}{d t}=Z \mathbf{z} \text { where } Z=Z_{2 n \times 2 n}=\left[\begin{array}{cc}
A & -I^{-2} B R^{-1} B^{H} \\
-Q & -A^{H}
\end{array}\right] \\
\mathbf{z}=\left[\begin{array}{l}
\mathbf{x} \\
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\end{array}\right], \quad \text { and } \quad\left\{\begin{array}{lll}
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( $Z$ has a Hamiltonian symmetry, such that eigenvalues appear in pairs of equal imaginary and opposite real part.)

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$$

(Z has a Hamiltonian symmetry, such that eigenvalues appear in pairs of equal imaginary and opposite real part.)

This linear ODE is a two-point boundary value problem and may be solved using a linear relationship between the state vector $\mathbf{x}(t)$ and adjoint vector $\mathbf{p}(t)$ vi a matrix $X(T)$ such that $\mathbf{p}=X \mathbf{x}$, and inserting this solution ansatz into (1) to eliminate $\mathbf{p}$.

## The Riccati equation

It follows that matrix $X$ obeys the differential Riccati equation

$$
\begin{equation*}
-\frac{d X}{d t}=A^{H} X+X A-X I^{-2} B R^{-1} B^{H} X+Q \quad \text { with } \quad X(T)=0 . \tag{2}
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Once $X$ is known, the optimal value of $\mathbf{u}$ may then be written in the form of a feedback control rule such that

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Finally, if the system is time invariant (LTI) and we take the limit that $T \rightarrow \infty$, the matrix X in (2) may be marched to steady state. This steady state solution for $X$ satisfies the continuous-time algebraic Riccati equation

$$
0=A^{H} X+X A-X I^{-2} B R^{-1} B^{H} X+Q,
$$

where additionally $X$ is constrained such that $A+B K$ is stable.

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V_{11} & V_{12} \\
V_{21} & V_{22}
\end{array}\right] \quad \text { and } \quad \mathbf{z}=\left[\begin{array}{l}
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the solutions $\mathbf{z}$ that obey the boundary conditions at $t \rightarrow \infty$ are spanned by the first $n$ columns of $V$. The direct ( $\mathbf{x}$ ) and adjoint ( $\mathbf{p}$ ) parts of the these columns are related as $\mathbf{p}=X \mathrm{x}$, where

$$
\left[\mathbf{p}_{1}, \mathbf{p}_{2}, \cdots, \mathbf{p}_{n}\right]=X\left[\mathbf{x}_{1}, \mathbf{x}_{2}, \cdots, \mathbf{x}_{n}\right] \quad \rightarrow \quad X=V_{21} V_{11}^{-1}
$$

## Motivation

- Optimal control via application of modern control algorithms (Riccati equation) is intractable because of the very large number of degrees of freedom deriving from the discretization of the Navier-Stokes equations.


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- One common approach is to use reduced-order models (ROM).
- Here, we present two exact methods which do not rely on such modeling,

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\text { where } \quad J=\frac{1}{2} \int_{0}^{T}\left[\mathbf{x}^{H} Q \mathbf{x}+I^{2} \mathbf{u}^{H} R \mathbf{u}\right] d t
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(3) In the limit that $l^{2} \rightarrow \infty$, MCE: Minimal Control Energy
(2) For any value of $l^{2}$, more general, ADA: Adjoint of the Direct-Adjoint

## Minimal Control Energy

## Minimal-energy control feedback I

In the limit that $I^{2} \rightarrow \infty$ we consider

$$
J=\int_{0}^{T} \frac{1}{2}\left[I^{-2} \mathbf{x}^{H} Q \mathbf{x}+\mathbf{u}^{H} R \mathbf{u}\right]
$$

With this defintion the same derivation as before leads to

$$
\frac{d \mathbf{z}}{d t}=Z \mathbf{z} \quad \text { where } \quad Z=Z_{2 n \times 2 n}=\left[\begin{array}{cc}
A & -B R^{-1} B^{H} \\
-I^{-2} Q & -A^{H}
\end{array}\right]
$$

Taking the limit $I^{2} \rightarrow \infty$ we get

## Minimal-energy control feedback II

In the limit that $I^{2} \rightarrow \infty$ we consider

$$
J=\int_{0}^{T} \frac{1}{2}\left[I^{-2} \mathbf{x}^{H} Q \mathbf{x}+\mathbf{u}^{H} R \mathbf{u}\right]
$$

With this defintion the same derivation as before leads to

$$
\frac{d \mathrm{z}}{d t}=Z \mathbf{z} \text { where } Z=Z_{2 n \times 2 n}=\left[\begin{array}{cc}
A & -B R^{-1} B^{H} \\
0 & -A^{H}
\end{array}\right]
$$

$Z$ becomes block triangular. The direct and adjoint equations are

$$
\frac{\partial \mathbf{x}}{\partial t}=A \mathbf{x}+B \mathbf{u}, \quad \mathbf{u}=-R^{-1} B^{H} \mathbf{p}, \quad \frac{\partial \mathbf{p}}{\partial t}=-A^{H} \mathbf{p}+0
$$

## Minimal-energy control feedback III

The eigenvalues of this system is given by the union of the eigenvalues of $A$ and $-A^{H}$.


The eigenvalues of $(+)$ the discretized open-loop system, and (o) the closed-loop system $A+B K$ after minimal-energy control is applied.

## Minimal-energy control feedback IV

Here we know the eigenvalues and only need to compute

$$
X=V_{21} V_{11}^{-1}
$$

It can be shown that $X$ is only function of $V_{21} . K$ is finally given as a function of the unstable eigenvalues and corresponding left eigenvectors.

$$
K=-B^{H} T_{u} F^{-1} T_{u}^{H}
$$

where $F$ has elements

$$
f_{i j}=c_{i j} /\left(\lambda_{i}+\lambda_{j}^{*}\right)
$$

and

$$
C=T_{u}^{H} B B^{H} T_{u}
$$

$T_{u}$ is the matrix containing unstable left eigenvectors

## ADA: Adjoint of the Direct-Adjoint

## Riccati-less optimal control I

The aim is to compute the solution for $K$, which is independent of $x_{0}$ and time invariant. This can be solved using an iterative procedure to "try" different $\mathrm{x}_{0}$ (computationally expensive).

## ALTERNATIVELY

For a converged solution at $t=0$ we can write

$$
\mathbf{u}=K \mathbf{x}_{0}=-\frac{1}{1^{2}} R^{-1} B^{H} \mathbf{p}_{0} .
$$

This is a linear relation between the input $\mathbf{x}_{0}$ and output $\mathbf{u}$.


The input has a large dimension and the output a small dimension.

## Riccati-less optimal control II

Such a problem is efficiently solved using the adjoint equations.

The adjoint input has a small dimension and the output a large dimension.

$K$ is obtained from the solution of the adjoint of the direct-adjoint system.

## Adjoint of the Direct-Adjoint system I

Introduce the adjoint variables $\mathbf{x}^{+}$and $\mathbf{p}^{+}$and multiply with the direct-adjoint equations, then integrate in time from $t=0$ to $t=T$. Obs! here we consider that $\mathbf{u}$ has dimension $m=1$.

$$
\int_{0}^{T} \mathbf{x}^{+H}\left(\frac{\partial \mathbf{x}}{\partial t}-A \mathbf{x}+\frac{1}{1^{2}} B R^{-1} B^{H} \mathbf{p}\right) d t+\int_{0}^{T} \mathbf{p}^{+H}\left(\frac{\partial \mathbf{p}}{\partial t}+A^{H} \mathbf{p}+Q \mathbf{x}\right) d t=0 .
$$

## Adjoint of the Direct-Adjoint system II

Using integration by parts, and considering that both $R$ and $Q$ are symmetric, we obtain

$$
\begin{aligned}
-\int_{0}^{T} \mathbf{p}^{H}\left(\frac{\partial \mathbf{p}^{+}}{\partial t}-A \mathbf{p}^{+}-\right. & \left.\frac{1}{\rho^{2}} B R^{-1} B^{H} \mathbf{x}^{+}\right) d t-\int_{0}^{T} \mathbf{x}^{H}\left(\frac{\partial \mathbf{x}^{+}}{\partial t}+A^{H} \mathbf{x}^{+}-Q \mathbf{p}^{+}\right) d t \\
& +\left[\mathbf{p}^{H} \mathbf{p}^{+}\right]_{0}^{T}+\left[\mathbf{x}^{H} \mathbf{x}^{+}\right]_{0}^{T}=0
\end{aligned}
$$

If we now define the new adjoint equations as

## Adjoint of the Direct-Adjoint system III

Using integration by parts, and considering that both $R$ and $Q$ are symmetric, we obtain

$$
\begin{gathered}
-\int_{0}^{T} \mathbf{p}^{H}(\underbrace{\frac{\partial \mathbf{p}^{+}}{\partial t}-A \mathbf{p}^{+}-\frac{1}{1^{2}} B R^{-1} B^{H} \mathbf{x}^{+}}_{=0}) d t-\int_{0}^{T} \mathbf{x}^{H}(\underbrace{\frac{\partial \mathbf{x}^{+}}{\partial t}+A^{H} \mathbf{x}^{+}-Q \mathbf{p}^{+}}_{=0}) d t \\
+\left[\mathbf{p}^{H} \mathbf{p}^{+}\right]_{0}^{T}+\left[\mathbf{x}^{H} \mathbf{x}^{+}\right]_{0}^{T}=0
\end{gathered}
$$

If we now define the new adjoint equations as

$$
\begin{gathered}
\frac{\partial \mathbf{p}^{+}}{\partial t}=A \mathbf{p}^{+}+\frac{1}{l^{2}} B R^{-1} B^{H} \mathbf{x}^{+} \\
\frac{\partial \mathbf{x}^{+}}{\partial t}=-A^{H} \mathbf{x}^{+}+Q \mathbf{p}^{+}
\end{gathered}
$$

## Adjoint of the Direct-Adjoint system IV

with $\mathbf{x}^{+}(t=T)=0$ and $\mathbf{p}(t=T)=0$, the remaining terms are

$$
\mathbf{x}^{+H}(0) \mathbf{x}(0)+\mathbf{p}^{+H}(0) \mathbf{p}(0)=0 .
$$

Recall that the original linear relation was

$$
K \mathrm{x}_{0}=-\frac{1}{1^{2}} R^{-1} B^{H} \mathbf{p}_{0}
$$

- Choosing $\mathbf{p}^{+H}(t=0)$ as one row of $-\frac{1}{1^{2}} R^{-1} B^{H}(m=1)$
- we can identify one row of $K$ as $\mathbf{x}^{+H}(0) .(m=1)$


## Riccati-less optimal control: solution procedure

If we let $\mathbf{x}^{+} \rightarrow-\mathbf{p}$ and $\mathbf{p}^{+} \rightarrow \mathbf{x}$ we easily obtain the original (Direct-Adjoint) system. (self-adjoint)

Finally: solve the original linear system with new b.c.

$$
\begin{aligned}
& \frac{\partial \mathbf{x}}{\partial t}=A \mathbf{x}-\frac{1}{1^{2}} B R^{-1} B^{H} \mathbf{p} \quad \text { on } \quad 0<t<T, \quad \mathbf{x}^{H}(0) \quad \text { is one row of } \frac{1}{1^{2}} R^{-1} B^{H}, \\
& \frac{\partial \mathbf{p}}{\partial t}=-A^{H} \mathbf{p}-Q \mathbf{x} \quad \text { on } \quad 0<t<T, \quad \text { with } \quad \mathbf{p}(T)=0 .
\end{aligned}
$$

One row of $K$ is then given by $-\mathbf{p}^{H}(0) \quad\left(\right.$ since $\left.\mathbf{x}^{+}=-\mathbf{p}\right)$.

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$$
\begin{aligned}
& \frac{\partial \mathbf{x}}{\partial t}=A \mathbf{x}-\frac{1}{l^{2}} B R^{-1} B^{H} \mathbf{p} \quad \text { on } \quad 0<t<T, \quad \mathbf{x}^{H}(0) \quad \text { is one row of } \frac{1}{1^{2}} R^{-1} B^{H}, \\
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> IMPORTANT

Avoid solving $X_{n \times n}$

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One row of $K$ is then given by $-\mathbf{p}^{H}(0) \quad\left(\right.$ since $\left.\mathbf{x}^{+}=-\mathbf{p}\right)$.

## IMPORTANT

Avoid solving $X_{n \times n} \quad$ solve original system $x_{n \times 1} m$ times

## Applications

## Applications

## MCE: Minimal Control Energy

In this method the feedback matrix $K$ is evaluated from the unstable open-loop solutions of the system.
Case: control of the cylinder wake (globally unstable flow)
Refs: Carini, Pralits, Luchini, JFS, 2013

## ADA: Adjoint of the Direct-Adjoint

This method is more general and does not depend on whether the system is unstable or not.
Cases: control of the cylinder wake, boundary layer transition
Refs: Pralits, Luchini, IUTAM Proceeding, 2010,
Semeraro, Pralits, Rowley, Henningson, JFM, 2013

## Control of the cylinder wake

## Control strategy

MCE \& ADA

## Numerical procedure

- All equations are discretized using second-order finite-differences over a staggered, stretched, Cartesian mesh.
- An immersed-boundary technique is used to enforce the boundary conditions on the cylinder.
- The nonlinear mean-flow equations, along with their boundary conditions, are solved by a Newton-Raphson procedure.
- The linear and nonlinear evolution equations are solved using Adams-Bashforth/Crank-Nicholson.
- The eigenvalue problems are solved using an Inverse Iteration algorithm
- Discrete adjoint equations (accurate to machine precision).


## Cases:

Reynolds numbers close to the first bifurcation, two-dimensional flow

## MCE

The linear feedback matrix $K$ which suppresses vortex shedding from a circular cylinder has been computed using:

## MCE

The linear feedback matrix $K$ which suppresses vortex shedding from a circular cylinder has been computed using:
Full state information,

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## MCE

The linear feedback matrix $K$ which suppresses vortex shedding from a circular cylinder has been computed using:
Full state information, Actuator: angular oscillation, $R e=U D / \nu$
Dimension of control $\mathbf{u}$ is $m=1$

$\square$

The feedback matrix $K(\mathbf{u}=K \mathbf{x})$

- $R e=55$
- $R e=75$
- $R e=100$
- $\operatorname{Re}=150$




## Results: linearized N-S equations

The temporal evolution of the frequency and growth rate is compared with the eigenvalue $\lambda$

- The Strouhal number: $S t=f D / U$ compared to $S t=\lambda_{r} / 2 \pi$
- The growth rate: $\sigma=\frac{d}{d t} \log (u(t))$ compared to $\lambda_{i}$

Test case: $\operatorname{Re}=55$, control is turned on at $t=18$



## Control of vortex shedding: $R e=55$

$\square$

## Control of vortex shedding: $R e=55$

$\square$

## Stationary vs. mean flow

St for limit cycle coincide with mean-flow eigenfrequency



## $K_{u}$ stationary vs. mean

- $R e=55$
- $R e=75$
- $R e=100$
- $R e=150$
$K_{u}$ stationary

$K_{u}$ mean



## $K_{v}$ stationary vs. mean

- $R e=55$
- $R e=75$
- $R e=100$
- $R e=150$
$K_{v}$ stationary

$K_{v}$ mean



## Control of vortex shedding: stationary vs. mean



Red: stationary flow, Green: mean flow

## ADA

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Full state information, Actuator: angular oscillation, $R e=U D / \nu$
Dimension of control $\mathbf{u}$ is $m=1$

$\square$

## Results: $K$ for $R e=55$

$$
K_{u}, I^{2}=1
$$


$K_{v}, I^{2}=1$


## Control of vortex shedding

In the temporal evolution of the lift $\left(C_{L}\right)$ and control $\mathbf{u}$ :

- $C_{L}$ and $\mathbf{u}$ tend to zero as the control is applied
- Control $\mathbf{u}$ strengthens as $I^{2}$ decrease

Test case: $R e=55$, control is turned on at $t=0$



## Control of vortex shedding

In the temporal evolution of drag $\left(C_{D}\right)$ coefficient:

- As the control is applied $C_{D}$ tends to the constant value corresponding to the steady state solution
- The control acts more quickly as $I^{2}$ is decreased

Test case: $R e=55$, control is turned on at $t=0$


## Control of the flat plate boundary layer I

## Linear quadratic controller (LQR)



$$
\begin{aligned}
\frac{d \mathbf{q}}{d t} & =\mathbf{A q}+\mathbf{B}_{2} \mathbf{u} \\
\mathbf{z} & =\left[\begin{array}{c}
\mathbf{C}_{1} \\
\mathbf{0}
\end{array}\right] \mathbf{q}+\left[\begin{array}{l}
0 \\
l
\end{array}\right] \mathbf{u}
\end{aligned}
$$

$$
\begin{aligned}
& \text { B2 - actuator } \\
& \text { C1 - objective function }
\end{aligned}
$$

$$
\mathcal{J}=\frac{1}{2} \mathbf{z}^{H} \mathbf{z}=\frac{1}{2} \int_{0}^{T}\left(\mathbf{q}^{H} \mathbf{C}_{1}^{H} \mathbf{C}_{1} \mathbf{q}+\mathbf{u}^{H} \mathbf{R} \mathbf{u}\right) d t
$$

Cost function - to be minimised

## Control of the flat plate boundary layer II

## Full dimensional gains



Gain L, x $=300$
Sensor upstream

Gain $L, x=410$
Sensor downstream

Gain K, $x=400$

Streamwise component
Estimation gain located upstream of the sensor (forward solution, AAD)
Control gain located downstream of the actuator (adjoint solution, ADA)

## Control of the flat plate boundary layer III



Semeraro, Pralits, Rowley, Henningson, JFM, 2013

## Some numerical issues

## Continuous vs. Discrete Adjoint Equations

The adjoint equations can be derived using two different approaches.
Both with advantages and disadvantages.
By definition we have

$$
\langle p, L x\rangle=
$$

- Continuous approach $\rightarrow$ : The adjoint eq derived by definition using the continuou: equations.
+ Straightforward derivation, reuse old cod $\epsilon$ programming
- Accuracy depends on discretization, diffic boundary conditions
- Discrete approach $\rightarrow$ : The adjoint equations are derived from the discretized direct equations.
+ Accuracy can be achieved close to machine precision, and can be independent of discretization !!
- Tricky derivation, usually requires making a new code, or larger changes of an existing code.


Here "def" means definition of the adjoint operator.
In the top row it is on continuous form while in the bottom row it is on discrete form.

## Derivation of the adjoint equation I

Consider the following optimal control problem (ODE) where $\phi$ is the state and $g$ the control.

$$
\frac{d \phi(t)}{d t}=-A \phi(t)+B g(t), \quad \text { for } \quad 0 \leq t \leq T
$$

with initial condition

$$
\phi(0)=\phi_{0}
$$

We can now define an optimization problem in which the goal is to find an optimal $g(t)$ by minimizing the following objective function

$$
J=\frac{\gamma_{1}}{2}[\phi(T)-\Psi]^{2}+\frac{\gamma_{2}}{2} \int_{0}^{T} g(t)^{2} d t
$$

## Derivation of the adjoint equation II

## Continuous approach

We can solve this problem using an adjoint identity approach or by introducing Lagrange multipliers.

$$
\int_{0}^{T} a\left[\frac{d \phi}{d t}+A \phi-B g\right] d t=\int_{0}^{T}\left[-\frac{d a}{d t}+A^{*} a\right] \phi d t-\int_{0}^{T} a B g d t+a(T) \phi(T)-a(0) \phi(0)
$$

If we now define the adjoint equation as $-d a / d t=-A^{*} a$ with an arbitrary initial condition $a(T)$ then the identity reduces to

$$
\mathrm{LHS}=-\int_{0}^{T} a B g d t+a(T) \phi(T)-a(0) \phi(0)
$$

By definition the Left Hand Side is identically zero but this is exactly what must be checked numerically, i.e. error $=\mid$ LHS $\mid$.

## Derivation of the adjoint equation III

The gradient of $J$ w.r.t. $g$ can be derived considering the $J$ is nonlinear in $\phi$ and $g$. We linearise by $\phi \rightarrow \phi+\delta \phi, g \rightarrow g+\delta g$ and then write the linearised objective function as

$$
\gamma_{1}[\phi(T)-\Psi] \delta \phi(T)=\delta J-\gamma_{2} \int_{0}^{T} g \delta g d t
$$

If we choose $a(T)=\gamma_{1}[\phi(T)-\Psi]$ then the equation for $\delta J$ can be substituted into the expression for the adjoint identity. If you further define the adjoint equations, remember that $\delta \phi(0)=0$, then the final identity is written

$$
\delta J=\int_{0}^{T}\left[\gamma_{2} g+B^{*} a\right] \delta g d t
$$

The adjoint equations and gradient of $J$ w.r.t. $g$ are written

$$
-\frac{d a}{d t}+A^{*} a, a(T)=\gamma_{1}[\phi(T)-\Psi], \quad \text { and } \quad \nabla J_{g}=\gamma_{2} g+B^{*} a
$$

The so called optimality condition is given by $\nabla J_{g}=0$.

## Derivation of the adjoint equation IV

- The accuracy of the adjoint solution is important since it quantfies a "gradient" in the optimization problem.
- The "error" must be evaluated to quantify the accuracy the adjoint solution.
- Note that the adjoint solution depends on the resolution $(\Delta t)$, and likewise the accuracy.
- Can we do better ?


## Derivation of the adjoint equation $V$

## Discrete approach

A discrete version of the direct equation is written
$\frac{\phi^{i+1}-\phi^{i}}{\Delta t}=-A \phi^{i}+B g^{i}, \quad$ for $\quad i=1, \ldots, N-1$,
where $N$ denotes the number of discrete points on the interval $[0, T], \Delta t$ is the constant time step, and

$$
\phi^{1}=\phi_{0},
$$

is the initial condition. This can be written as a discrete evolution equation
$\phi^{i+1}=[I-\Delta t A] \phi^{i}+\Delta t B g^{i}, \quad$ for $\quad i=1, \ldots, N-1$.
A discrete version of the objective function can be written

$$
J=\frac{\gamma_{1}}{2}\left(\phi^{N}-\Phi\right)^{2}+\frac{\gamma_{2}}{2} \sum_{i=1}^{N-1} \Delta t\left(g^{i}\right)^{2}
$$

$$
\text { 1. } a^{i}
$$

An adjoint variable $a^{i}$ is introduced defined on $i=1, \ldots, N$ and by definition
$a^{i+1} \cdot L \phi^{i}=\left(L^{\star} a^{i+1}\right) \cdot \phi^{i}, \quad$ for $\quad i=1, \ldots, N-1$.
We then introduce the definition of the state equation on the left hand side of and impose that

$$
a^{i}=L^{\star} a^{i+1} \quad \text { for } \quad i=N-1, \ldots, 1
$$

This is the discrete adjoint equation. Using the discrete direct and adjoint yields

$$
a^{i+1} \cdot\left(\phi^{i+1}-\Delta t B g^{i}\right)=a^{i} \cdot \phi^{i}, \quad \text { for } \quad i=1, \ldots, N-1
$$ which must be valid for any $\phi$ and $a$. An error can therefore be written as

$$
\text { error }=\left|a^{N} \cdot \phi^{N}-a^{1} \cdot \phi^{1}-\sum_{i=1}^{N-1} \Delta t a^{i+1} \cdot B g^{i}\right|
$$

## Derivation of the adjoint equation VI

The discrete optimality condition is then derived. Since $J$ is nonlinear with respect to $\phi$ and $g$ we must first linearize. This can be written

$$
\delta J=\gamma_{1}\left(\phi^{N}-\Phi\right) \cdot \delta \phi^{N}+\gamma_{2} \sum_{i=1}^{N-1} \Delta t g^{i} \cdot \delta g^{i}
$$

We now choose the terminal condition of the adjoint as $a^{N}=\gamma_{1}\left(\phi^{N}-\Phi\right)$ and substitute this expression into the discrete adjoint identity. This is written

$$
\gamma_{1}\left(\phi^{N}-\Phi\right) \cdot \delta \phi^{N}=a^{1} \cdot \delta \phi^{1}+\sum_{i=1}^{N-1} \Delta t a^{i+1} \cdot B \delta g^{i}
$$

By inspection one can see that the left hand side is identical to the first term in the expression for $\delta J$, and $\delta \phi^{1}=0$. Rearranging the terms, we get

$$
\delta J=\sum_{i=1}^{N-1} \Delta t\left(\gamma_{2} g^{i}+B^{\star} a^{i+1}\right) \cdot \delta g^{i}
$$

from which we get the discrete optimality condition

$$
g^{i}=-\frac{1}{\gamma_{2}} B^{\star} a^{i+1} \quad \text { for } \quad i=1, . ., N-1
$$

Note that if $B$ is a matrix then $B^{\star}=B^{T}$.

## Checkpointing algorithm

- When the adjoint equation is forced in time by the direct solution (ex. quadratic objective function), then this poses storage requirements (hard ware). This becomes a problem for 2D and 3D problems with high resolution in space and time.
- One way to come around this is to apply Checkpointing. This consists of sampling the direct solution at given rate and then recompute the direct solution for short time intervals when needed. This means in theory that one more solution of the direct system has been added to the computational effort.
- However, since it is common to use parallel computing, and processors is becoming a smaller problem on can do something to obtain the minimal required computational time.
- This is done by rec :omputing the adjoint.



## EXTRA SLIDES

## Background: control using rotational oscillation

Aim: reduce $C_{D}$
Exp. Tokumaru \& Dimotakis (1991), $-20 \%, \operatorname{Re}=15000$
Feedback control:
Exp. Fujisawa \& Nakabayashi (2002) -16\% (-70\% C $C_{L}$ ), $R e=20000$
Exp. Fujisawa et al.(2001) "reduction", $R e=6700$
Optimal control (using adjoints):
Num. He et al.(2000) -30 to -60\% for $R e=200-1000$
Num. Protas \& Styczek (2002) -7\% at $R e=75,-15 \%$ at $R e=150$ Bergmann et al.(2005) $-25 \%$ at $R e=200$ (POD)

Aim: reduce vortex shedding Feedback control:
Num. Protas (2004) reduction, "point vortex model", $R e=75$
Optimal control (using adjoints):
Num. Homescu et al.(2002) reduction, $R e=60-1000$

## Minimal-energy control feedback

Denoting:

- $\mathbf{x}^{i}$ and $\lambda^{i}$ the $i$-th right eigenvector and eigenvalue of $A$,
- $\mathbf{y}^{i}$ and $-\lambda^{i *}$ the $i$-th right eigenvector and eigenvalue of $-A^{H}$,
- $\mathbf{y}^{i *}$ is left eigenvector of $A$, we see that the stable eigenvectors of

$$
\frac{\partial \mathbf{x}}{\partial t}=A \mathbf{x}+B \mathbf{u}, \quad \mathbf{u}=-R^{-1} B^{H} \mathbf{p}, \quad \frac{\partial \mathbf{p}}{\partial t}=-A^{H} \mathbf{p}
$$

are of two possible types:

$$
\begin{array}{lll}
\mathbf{p}=0, \mathbf{x}=\mathbf{x}^{i} & \text { if } & \Re\left(\lambda^{i}\right)<0 \\
\mathbf{p}=\mathbf{y}^{i}, \mathbf{x}=\left(\lambda^{i *}+A\right)^{-1} B R^{-1} B^{H} \mathbf{y}^{i} & \text { if } & \Re\left(\lambda^{i}\right)>0
\end{array} \text { (unstable) }
$$

We now project an arbitrary initial condition $\mathrm{x}_{0}$ onto these modes,

$$
\begin{equation*}
\mathbf{x}_{0}=\sum_{\text {stable }} d_{j} x^{j}+\sum_{\text {unstable }} f_{j}\left(\lambda^{j *}+A\right)^{-1} B R^{-1} B^{H} y^{j} \tag{4}
\end{equation*}
$$

and note that in order to reconstruct $\mathbf{p}$ we only need the $f_{j}$ 's, because the stable modes have $\mathbf{p}=0$. The coefficients $d_{j}$ can be eliminated from (4) by projecting the left eigenvectors:

$$
\mathbf{y}^{i *} \mathbf{x}_{0}=\mathbf{y}^{i *} \sum_{\text {unstable }} f_{j}\left(\lambda^{j *}+A\right)^{-1} B R^{-1} B^{H} \mathbf{y}^{j}=\sum_{\text {unstable }} c_{i j} f_{j}
$$

where, since $\boldsymbol{y}^{i *}$ is also a left eigenvector of $\left(\lambda^{j *}+A\right)^{-1}$,

$$
c_{i j}=\frac{\mathbf{y}^{i *} B R^{-1} B^{H} \mathbf{y}^{j}}{\lambda^{i}+\lambda^{j *}}
$$

Only the unstable eigenvalues and left eigenvectors are needed.

## The main theorem

Summarizing, the solution of the minimal-energy stabilizing control feedback problem can be written in terms of the unstable left eigenvectors only.

Theorem 1. Consider a stabilizable system $\dot{x}=A \mathbf{x}+$ Bu with no pure imaginary open-loop eigenvalues. Determine the unstable eigenvalues and corresponding left eigenvectors of $A$ such that $T_{u}^{H} A=\Lambda_{u} T_{u}^{H}$ (equivalently, determine the unstable eigenvalues and corresponding right eigenvectors of $A^{H}$ such that $\left.A^{H} T_{u}=T_{u} \Lambda_{u}^{H}\right)$. Define $\bar{B}_{u}=T_{u}^{H} B$ and $C=\bar{B}_{u} \bar{B}_{u}^{H}$, and compute a matrix $F$ with elements $f_{i j}=c_{i j} /\left(\lambda_{i}+\lambda_{j}^{*}\right)$. The minimal-energy stabilizing feedback controller is then given by $\mathbf{u}=K \mathbf{x}$, where $K=-\bar{B}_{u}^{H} F^{-1} T_{u}^{H}$.

